

HEATING RATES

DEFINITIONS

Prandtl Number

$$Pr = \mu \frac{c_p}{K}$$

where $\left\{ \begin{array}{l} c_p = \text{specific heat @} \\ \text{constant pressure} \\ K = \text{thermal conductivity} \\ \mu = \text{viscosity} \end{array} \right\}$

$$Pr \propto \frac{\text{frictional dissipation}}{\text{thermal conduction}}$$

$$Pr \approx 0.715 \text{ for air at standard conditions}$$

Sutherland's Law (empirical)

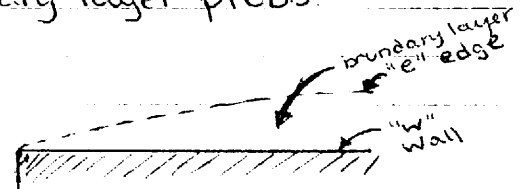
viscosity depends on temp

$$\frac{\mu}{\mu_{ref}} = \left(\frac{T}{T_{ref}} \right)^{3/2} \frac{T_{ref} + S}{T + S}$$

$$\text{for air: } \left\{ \begin{array}{l} \mu_{ref} = 1.789 \times 10^{-5} \text{ Kg/m.s} \\ T_{ref} = 288 \text{ K} \\ S = 110 \text{ K} \\ \cdot \text{good up to several thousand degrees} \end{array} \right.$$

Stanton Number - applies to boundary layer probs.

$$St = \frac{\dot{q}_w}{\rho_e v_e (H_{aw} - H_w)}$$



"e" means edge of boundary layer
"w" means wall

$H = \text{enthalpy}$
 $= c_p T$ — for perfect gas

$H_w = \text{enthalpy @ the wall}$

$H_{aw} = \text{enthalpy @ the wall if it were adiabatic}$

$$\rightarrow \text{for } H_{aw}, \left. \frac{\partial T}{\partial z} \right|_{\text{wall}} = 0$$

Approximating H_{aw}

$$H_{aw} = H_e + r \frac{u_e^2}{2}$$

, r = recovery factor

$$H_o = H_e + \frac{u_e^2}{2}$$

← TOTAL ENTHALPY at edge of boundary layer

so... $H_{aw} = H_e + r (H_o - H_e)$

for incompressible flow, $r \approx \sqrt{Pr}$
= 0.845 for std. air

r decreases only 2.4% from $M=0$ to $M=16$
→ fairly constant!

Reynold's Analogy - pg 81 in Hanky

$$\frac{St}{C_f} = \frac{1}{2} Pr^{-2/3}$$

, C_f = skin friction coeff.

since $Pr \approx 1$,

$$St \approx \frac{C_f}{2}$$

this is the Reynold's Analogy

Empirical Correlation for C_f :

$$\frac{C_f}{2} = \frac{A}{Re^n}$$

	A	n
LAMINAR FLOW	0.332	0.5
TURBULENT FLOW	0.0296	0.2

transition btwn laminar & turbulent @ $\sim Re = 10^6$

$$\rightarrow St \approx \frac{A}{Re^n} Pr^{-2/3}$$

STANTON NUMBER

FRACTION OF HEAT INTO WAKE

$$\dot{q}_A = \rho A V H_s \quad \left. \begin{array}{l} \uparrow \\ \text{total enthalpy} \end{array} \right\}$$

thermal energy flow rate
in free stream

$$\dot{q}_w = \dot{q}_B = S_T \rho A V (H_{aw} - H_w) \quad \left. \right\}$$

amount absorbed by vehicle
(EON 4.3, pg 81)

$$h = \left[k \frac{\partial T}{\partial y} \right]_{\text{wall}}$$

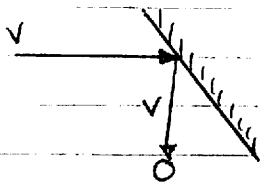
usually, $H_{aw} \gg H_w$
 $H_{aw} \approx H_s$

$$\dot{q}_B = S_T \rho A V H_s$$

$$\frac{\dot{q}_B}{\dot{q}_A} = S_T \approx 10^{-3}$$

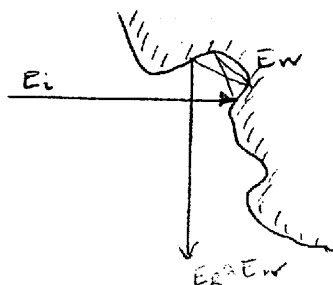
→ 0.1% of heating goes into vehicle!
most of the energy goes into the atmosphere

FREE MOLECULAR FLOW



at equilibrium
- no heating

on a larger scale...

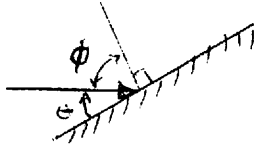


more accurately
(rough surface, molecule
bounces a few times
before leaving vehicle)

Newtonian Flow:

$$C_p = Z \sin^2 \theta = Z \cos^2 \phi$$

$\phi = \text{complement of } \theta$



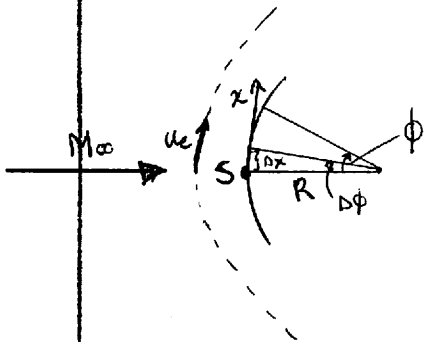
$$\frac{P_e - P_\infty}{q_\infty} = Z \cos^2 \phi$$

$$P_e = Z q_\infty \cos^2 \phi + P_\infty$$

$$\frac{dP_e}{dx} = -4 q_\infty \cos \phi \sin \phi \frac{d\phi}{dx}$$

$$\frac{dU_e}{dx} = \frac{4 q_\infty}{\rho_e U_e} \cos \phi \sin \phi \frac{d\phi}{dx} \quad \text{everywhere on the body}$$

Now Evaluate Heating Rate at STAGNATION POINT, S



\approx incompressible flow around S

$$U_e = \left(\frac{dU_e}{dx} \right)_s \Delta x$$

$$\text{at } s: \quad \cos \Delta \phi \approx 1$$

$$\sin \Delta \phi \approx \Delta \phi = \frac{\Delta x}{R}$$

$$\frac{\Delta \phi}{\Delta x} = \frac{1}{R}$$

$$C_p = Z \cos^2 \Delta \phi \approx Z$$

$$q_\infty = \frac{1}{2} (\rho_e U_e^2)$$

$$\left(\frac{dU_e}{dx} \right)_s^2 = \frac{Z (\rho_e U_e^2)}{\rho_e \Delta x} \left(\frac{\Delta x}{R} \right) \left(\frac{1}{R} \right) \rightarrow \frac{dU_e}{dx} = \frac{1}{R} \sqrt{\frac{Z (\rho_e U_e^2)}{\rho_e}}$$

Small angle approx
 $\sin \theta \approx \theta$
 $\cos \theta \approx 1$

$q_w \propto \frac{1}{\sqrt{R}}$, BLUNT BODIES STAY COOLER THAN SHARP ONES

Got to
here
2/24/98

APPROXIMATIONS FOR HEATING RATES

(these will get you within 10% or so)

Simple Set of Equations for HYPERSONIC CONTINUUM FLOW

$$\dot{q}_w = \rho_{\infty}^N V_{\infty}^M C$$

$$\begin{array}{ll} \dot{q}_w & \text{in } \text{W/cm}^2 \\ \rho_{\infty} & \text{in } \text{kg/m}^3 \\ V_{\infty} & \text{in } \text{m/s} \\ R & \text{in } \text{m} \end{array}$$

Stagnation Point:

$$M = 3$$

$$N = 0.5$$

$$C = 1.83 \times 10^{-8} R^{-1/2} \left(1 - \frac{H_w}{H_o}\right)$$

Laminar Flow on Flat Plate:

$$M = 3.2$$

$$N = 0.5$$

$$C = 2.53 \times 10^{-9} (\cos \phi)^{1/2} (\sin \phi) x^{-1/2} \left(1 - \frac{H_w}{H_o}\right)$$

where x = distance along surface in m

Turbulent Flow on Flat Plate:

CASE I: $V_{\infty} \leq 3962 \text{ m/s}$

$$M = 3.37$$

$$N = 0.8$$

$$C = 3.9 \times 10^{-8} (\cos \phi)^{1.78} (\sin \phi)^{1.6} x_T^{-1/5} \left(\frac{T_w}{556}\right)^{-1/4} \left(1 - 1.1 \frac{H_w}{H_o}\right)$$

CASE II: $V_{\infty} > 3962 \text{ m/s}$

$$M = 3.7$$

$$N = 0.8$$

$$C = 3.2 \times 10^{-7} (\cos \phi)^{2.08} (\sin \phi)^{1.6} x_T^{-1/5} \left(1 - 1.1 \frac{H_w}{H_o}\right)$$

where x_T = distance measured in the turbulent part of the boundary layer