

# Ballistic Atmospheric Entry

- Standard atmospheres
- Orbital decay due to atmospheric drag
- Straight-line (no gravity) ballistic entry based on atmospheric density
- Straight-line (no gravity) ballistic entry based on altitude, rather than density
- Planetary entries (at least a start)

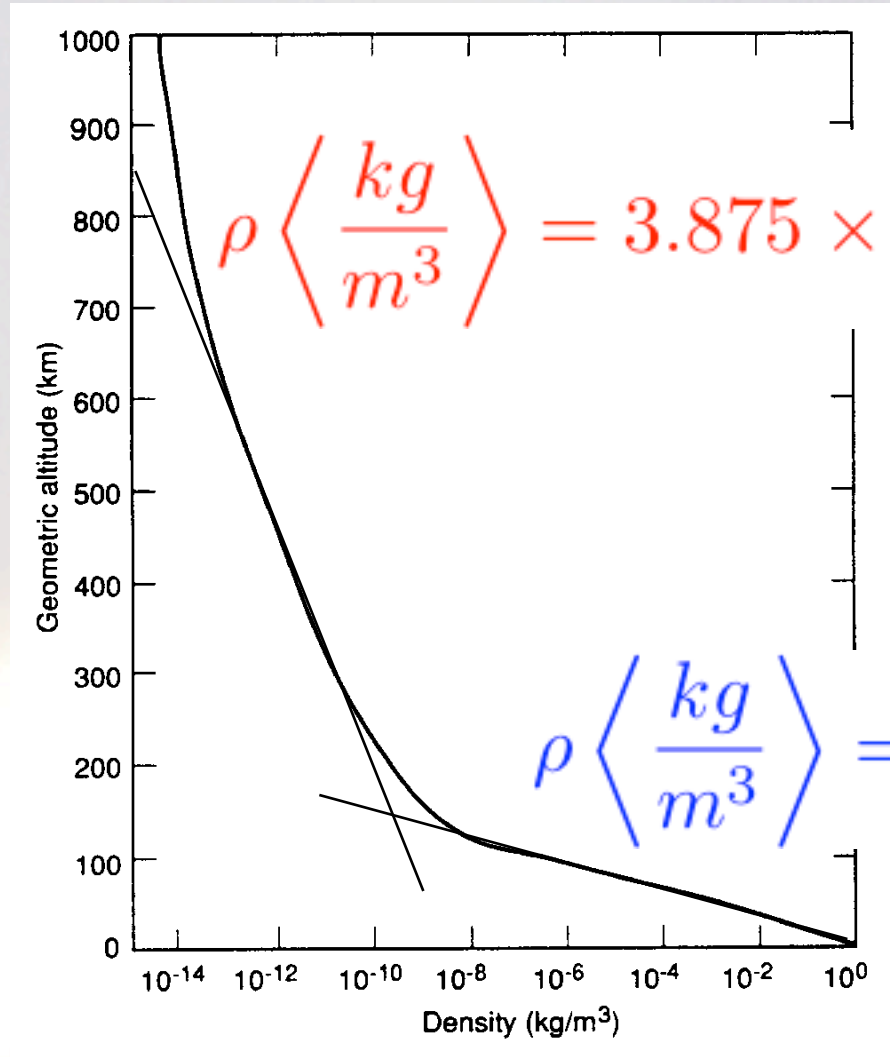
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UNIVERSITY OF  
MARYLAND

**Ballistic Atmospheric Entry II**  
**ENAE 791 - Launch and Entry Vehicle Design**

# Atmospheric Density with Altitude



$$\rho \left\langle \frac{kg}{m^3} \right\rangle = 3.875 \times 10^{-9} e^{-\frac{h \langle km \rangle}{59.06}}$$

$$\rho \left\langle \frac{kg}{m^3} \right\rangle = 1.226 e^{-\frac{h \langle km \rangle}{7.524}}$$

Ref: V. L. Pisacane and R. C. Moore, Fundamentals of Space Systems Oxford University Press, 1994



# Energy Loss Due to Atmospheric Drag

$$\text{Drag } D \equiv \frac{1}{2} \rho v^2 A c_D$$

$$\text{Drag acceleration } a_d = \frac{D}{m} = \frac{\rho v^2}{2} \frac{A c_D}{m}$$

$$\beta \equiv \frac{m}{c_D A}$$

<== Ballistic Coefficient

$$a_d = \frac{\rho v^2}{2\beta}$$

$$\text{orbital energy } \equiv E = -\frac{\mu}{2a}$$

$$\frac{dE}{dt} = \frac{\mu}{2a^2} \frac{da}{dt}$$



# Energy Loss Due to Atmospheric Drag

Since drag is highest at perigee, the first effect of atmospheric drag is to circularize the orbit (high perigee drag lowers apogee)

$$\frac{dE_{drag}}{dt} = a_d v$$

$$v_{circ}^2 = \frac{\mu}{a} \quad \frac{dE_{drag}}{dt} = -\frac{\rho v^2}{2\beta} \sqrt{\frac{\mu}{a}}$$

$$\frac{dE_{drag}}{dt} = -\sqrt{\frac{\mu}{a}} \frac{\rho}{2\beta} \frac{\mu}{a} = -\left(\frac{\mu}{a}\right)^{\frac{3}{2}} \frac{\rho}{2\beta}$$



# Derivation of Orbital Decay Due to Drag

Set orbital energy variation equal to energy lost by drag

$$\frac{\mu}{2a^2} \frac{da}{dt} = -\frac{\rho}{2\beta} \left(\frac{\mu}{a}\right)^{\frac{3}{2}}$$

$$\frac{da}{dt} = -\frac{\rho}{\beta} \sqrt{\mu a}$$

$$\rho = \rho_0 e^{-\frac{h}{h_s}} \quad a = h + r_E \implies \frac{da}{dt} = \frac{dh}{dt}$$

$$\frac{dh}{dt} = -\frac{\sqrt{\mu (h + r_E)}}{\beta} \rho_0 e^{-\frac{h}{h_s}}$$



# Derivation of Orbital Decay (2)

This is a separable differential equation...

$$\frac{1}{\sqrt{r_E + h}} e^{\frac{h}{h_s}} dh = -\frac{\sqrt{\mu}}{\beta} \rho_o dt$$
$$\int_{h_o}^h \frac{1}{\sqrt{r_E + h}} e^{\frac{h}{h_s}} dh = -\frac{\sqrt{\mu}}{\beta} \rho_o \int_{t_o}^t dt$$

Assume  $\sqrt{r_E + h} \sim \sqrt{r_E}$  for  $r_E \gg h$

$$\frac{1}{\sqrt{r_E}} \int_{h_o}^h e^{\frac{h}{h_s}} dh = -\frac{\sqrt{\mu}}{\beta} \rho_o (t - t_o)$$



# Derivation of Orbital Decay (3)

$$\frac{h_s}{\sqrt{r_E}} \left( e^{\frac{h}{h_s}} - e^{\frac{h_o}{h_s}} \right) = -\frac{\sqrt{\mu}}{\beta} \rho_o (t - t_o)$$

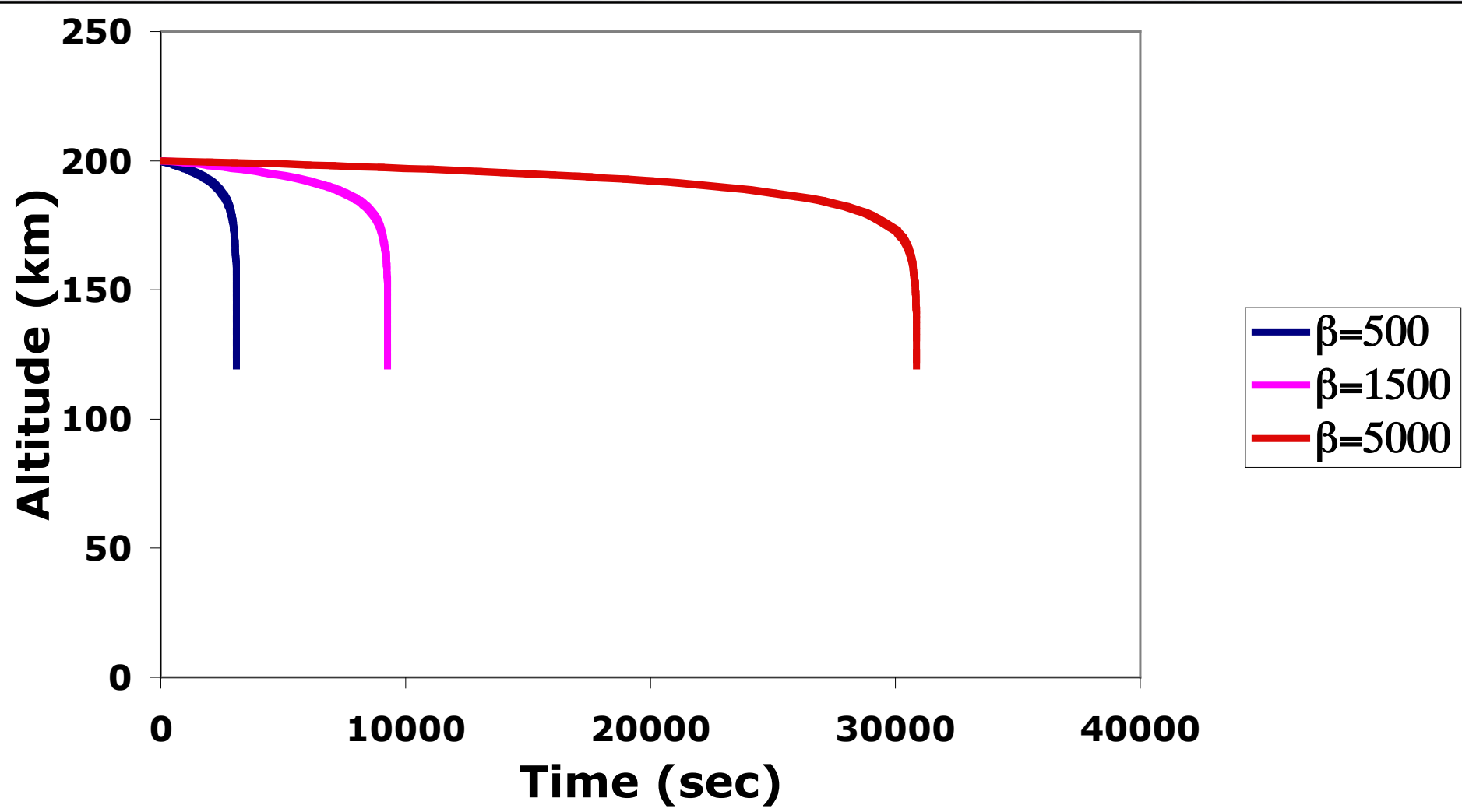
$$e^{\frac{h}{h_s}} - e^{\frac{h_o}{h_s}} = -\frac{\sqrt{\mu r_E}}{h_s \beta} \rho_o (t - t_o)$$

$$h(t) = h_s \ln \left[ e^{\frac{h_o}{h_s}} - \frac{\sqrt{\mu r_E}}{h_s \beta} \rho_o (t - t_o) \right]$$

Note that some variables typically use km, and others are in meters - you have to make sure unit conversions are done properly to make this work out correctly!



# Orbit Decay from Atmospheric Drag



# Time Until Orbital Decay

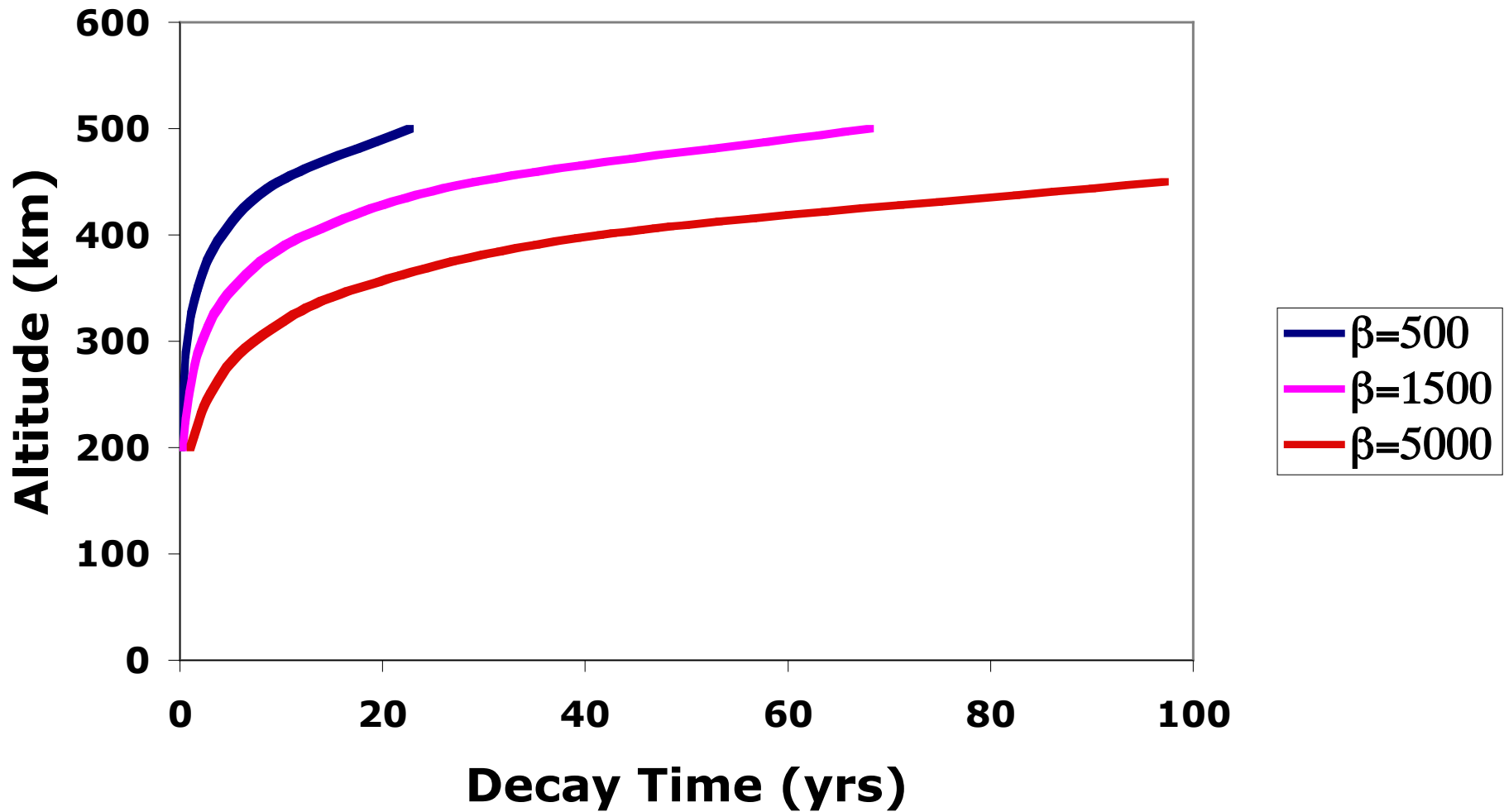
$$e^{\frac{h}{h_s}} - e^{\frac{h_o}{h_s}} = -\frac{\sqrt{\mu r_E}}{h_s \beta} \rho_o (t - t_o)$$

To find the time remaining ( $t_o=0$ ) until the orbit reaches any given “critical” altitude, some algebra gives

$$t(h) = \frac{h_s \beta}{\sqrt{\mu r_E} \rho_o} \left( e^{\frac{h_o}{h_s}} - e^{\frac{h}{h_s}} \right)$$



# Decay Time to $r=120$ km



# Ballistic Entry (no lift)

$s$  = distance along the flight path

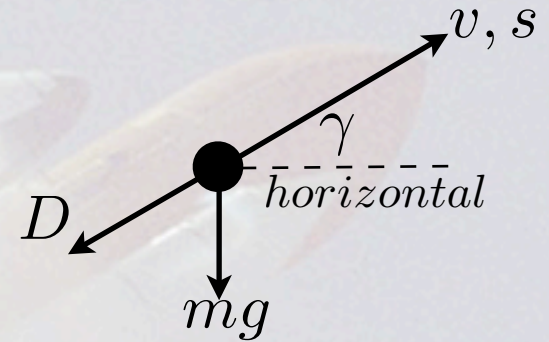
$$\frac{dv}{dt} = -g \sin \gamma - \frac{D}{m}$$

$$\frac{dv}{dt} = \frac{dv}{ds} \frac{ds}{dt} = V \frac{dv}{ds} = \frac{1}{2} \frac{d(v^2)}{ds}$$

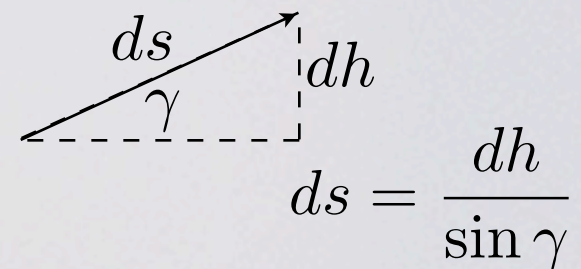
$$\frac{1}{2} \frac{d(v^2)}{ds} = -g \sin \gamma - \frac{D}{m}$$

$$\frac{1}{2} \frac{d(v^2)}{ds} = -g \sin \gamma - \frac{\rho v^2}{2m} A c_D$$

$$\frac{\sin \gamma}{2} \frac{d(v^2)}{dh} = -g \sin \gamma - \frac{\rho v^2}{2m} A c_D$$



$$\text{Drag } D \equiv \frac{1}{2} \rho v^2 A c_D$$



# Ballistic Entry (2)

Exponential atmosphere  $\Rightarrow \rho = \rho_o e^{-\frac{h}{h_s}}$

$$\frac{d\rho}{\rho_o} = e^{-\frac{h}{h_s}} \left( \frac{-dh}{h_s} \right) = \frac{\rho_o e^{-\frac{h}{h_s}}}{\rho_o} \left( \frac{-dh}{h_s} \right) = \frac{\rho}{\rho_o} \left( \frac{-dh}{h_s} \right)$$

$$dh = \frac{-h_s}{\rho} d\rho$$

$$\frac{\sin \gamma}{2} \frac{d(v^2)}{dh} = -g \sin \gamma - \frac{\rho v^2}{2m} A c_D$$

$$\frac{\sin \gamma}{2} \frac{d(v^2)}{d\rho} \left( \frac{-\rho}{h_s} \right) = -g \sin \gamma - \frac{\rho v^2}{2} \frac{A c_D}{m}$$

$$\frac{d(v^2)}{d\rho} = \frac{2gh_s}{\rho} + \frac{h_s v^2}{\sin \gamma} \frac{A c_D}{m}$$



# Ballistic Entry (3)

Let  $\beta \equiv \frac{m}{c_D A} \Rightarrow$  Ballistic Coefficient

$$\frac{d(v^2)}{d\rho} - \frac{h_s}{\beta \sin \gamma} v^2 = \frac{2gh_s}{\rho}$$

Assume  $mg \ll D$  to get homogeneous ODE

$$\frac{d(v^2)}{d\rho} - \frac{h_s}{\beta \sin \gamma} v^2 = 0 \qquad \frac{d(v^2)}{v^2} = \frac{h_s}{\beta \sin \gamma} d\rho$$

Use  $(v^2)$  as integration variable

$$\int_{v_e}^v \frac{d(v^2)}{v^2} = \frac{h_s}{\beta \sin \gamma} \int_0^\rho d\rho \qquad v_e = \text{velocity at entry}$$



# Ballistic Entry (4)

Note that the effect of ignoring gravity is that there is no force perpendicular to velocity vector  $\Rightarrow$  constant flight path angle  $\gamma$   
 $\Rightarrow$  straight line trajectories

$$\ln \frac{v^2}{v_e^2} = 2 \ln \frac{v}{v_e} = \frac{h_s \rho}{\beta \sin \gamma}$$

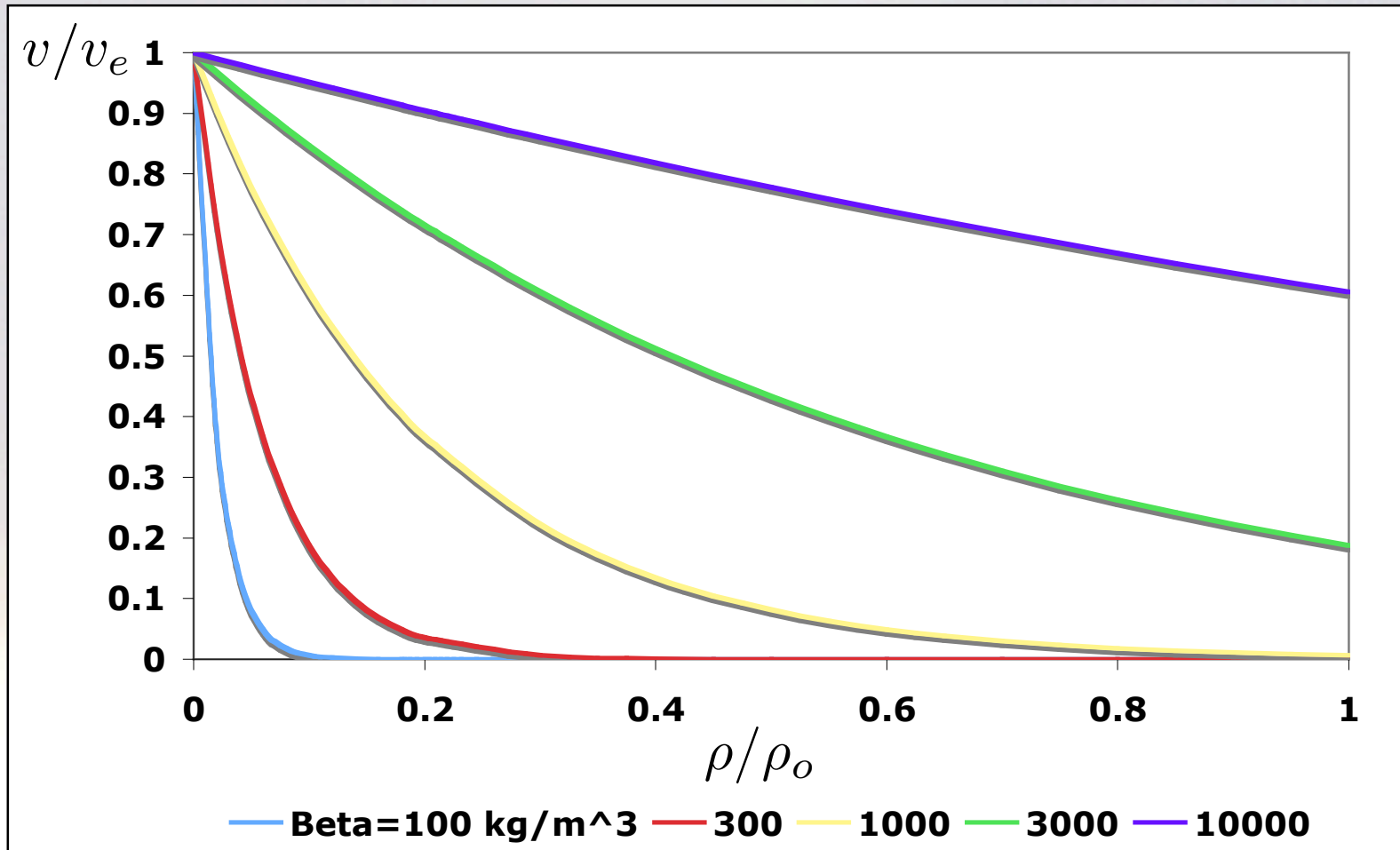
$$\frac{v}{v_e} = \exp \left( \frac{h_s \rho}{2\beta \sin \gamma} \right)$$

$$\frac{v}{v_e} = \exp \left( \frac{h_s \rho_o}{2\beta \sin \gamma} \frac{\rho}{\rho_o} \right)$$

Check units:  $\left( \frac{m \frac{kg}{m^3}}{\frac{kg}{m^2}} \right)$



# Earth Entry, $\gamma = -60^\circ$



# What About Peak Deceleration?

$$n \equiv \frac{dv}{dt} = -\frac{\rho v^2}{2\beta}$$

To find  $n_{max}$ , set  $\frac{d}{dt} \left( \frac{dv}{dt} \right) = \frac{d^2v}{dt^2} = 0$

$$\frac{d^2v}{dt^2} = -\frac{1}{2\beta} \left( 2\rho v \frac{dv}{dt} + v^2 \frac{d\rho}{dt} \right) = 0$$

$$\frac{d^2v}{dt^2} = -\frac{1}{2\beta} \left( -\frac{2\rho^2 v^3}{2\beta} + v^2 \frac{d\rho}{dt} \right) = 0$$

$$\frac{\rho^2 v^3}{\beta} = v^2 \frac{d\rho}{dt} \qquad \rho^2 v = \beta \frac{d\rho}{dt}$$



# Peak Deceleration (2)

From exponential atmosphere,

$$\frac{d\rho}{dt} = -\frac{\rho_o}{h_s} e^{-\frac{h}{h_s}} \frac{dh}{dt} = -\frac{\rho}{h_s} \frac{dh}{dt}$$

From geometry,  $\frac{dh}{dt} = v \sin \gamma$

$$\frac{d\rho}{dt} = -\frac{\rho v}{h_s} \sin \gamma \quad \rho^2 v = \beta \frac{d\rho}{dt}$$

$$\rho^2 v = \beta \left( -\frac{\rho v}{h_s} \sin \gamma \right)$$

Remember that this refers to the conditions at max deceleration

$$\rho n_{max} = -\frac{\beta}{h_s} \sin \gamma$$



# Critical $\beta$ for Deceleration Before Impact

At surface,  $\rho = \rho_o$

$$\beta_{crit} = -\frac{\rho_o h_s}{\sin \gamma} \quad \Leftarrow \text{Value of } \beta \text{ at which vehicle hits ground at point of maximum deceleration}$$

How large is maximum deceleration?

$$\frac{dv}{dt} = \frac{\rho v^2}{2\beta} \quad \Rightarrow \quad \left| \frac{dv}{dt} \right|_{max} = \frac{\rho_{n_{max}} v^2}{2\beta}$$

$$\left| \frac{dv}{dt} \right|_{max} = \frac{v^2}{2\beta} \left( -\frac{\beta}{h_s} \sin \gamma \right) = -\frac{1}{2} \frac{v^2}{h_s} \sin \gamma$$

Note that this value of  $v$  is actually  $v_{n_{max}}$



# Peak Deceleration (3)

From page 14,

$$\frac{v}{v_e} = \exp\left(\frac{h_s \rho}{2\beta \sin \gamma}\right)$$

$$\frac{v_{n_{max}}}{v_e} = \exp\left[\frac{h_s}{2\beta \sin \gamma} \left(-\frac{\beta}{h_s} \sin \gamma\right)\right] = e^{-\frac{1}{2}}$$

$$\left|\frac{dv}{dt}\right|_{max} = -\frac{1}{2} \frac{\left(v_e e^{-\frac{1}{2}}\right)^2}{h_s} \sin \gamma = -\frac{v_e^2 \sin \gamma}{2h_s e}$$

Note that the velocity at which maximum deceleration occurs is always a fixed fraction of the entry velocity - it doesn't depend on ballistic coefficient, flight path angle, or anything else! Also, the magnitude of the maximum deceleration is not a function of ballistic coefficient - it is dependent on the entry trajectory ( $v_e$  and  $\gamma$ ) but not spacecraft parameters (i.e., ballistic coefficient).



# Terminal Velocity

Full form of ODE -

$$\frac{d(v^2)}{d\rho} - \frac{h_s}{\beta \sin \gamma} v^2 = \frac{2gh_s}{\rho}$$

At terminal velocity,  $v = \text{constant} \equiv v_T$

$$-\frac{h_s}{\beta \sin \gamma} v_T^2 = \frac{2gh_s}{\rho}$$

$$v_T^2 = \sqrt{-\frac{2g\beta \sin \gamma}{\rho}}$$



# “Cannon Ball” $\gamma = -90^\circ$ Ballistic Entry

6.75” diameter sphere,  $c_D = 0.2$ ,  $V_E = 6000$  m/sec

	Iron	Aluminum	Balsa Wood
<b>Weight</b>	<b>40 lb</b>	<b>15.6 lb</b>	<b>14.5 oz</b>
$\beta_{md}$ (kg/m <sup>2</sup> )	<b>3938</b>	<b>1532</b>	<b>89</b>
$\rho_{md}$ (kg/m <sup>3</sup> )	<b>0.555</b>	<b>0.216</b>	<b>0.0125</b>
$h_{md}$ (m)	<b>5600</b>	<b>12,300</b>	<b>32,500</b>
$V_{impact}$ (m/s)	<b>1998</b>	<b>355</b>	<b>0*</b>
$V_{term}$ (m/sec)	<b>251</b>	<b>156</b>	<b>38</b>

\*Artifact of assumption that  $D \gg mg$



# Atmospheric Density with Altitude

Pressure=the integral of the atmospheric density in the column above the reference area

$$\rho = f(h) \quad P_o = \int_0^{\infty} \rho g dh = \rho_o g \int_0^{\infty} e^{-\frac{h}{h_s}} dh = -\rho_o g h_s \left[ e^{-\frac{h}{h_s}} \right]_0^{\infty} \\ = -\rho_o g h_s [0 - 1]$$

$$P_o = \rho_o g h_s$$

$$\text{Earth: } \rho_o = 1.226 \frac{\text{kg}}{\text{m}^3}; h_s = 7524\text{m};$$

$$P_o(\text{calc}) = 90,400 \text{ Pa}; P_o(\text{act}) = 101,300 \text{ Pa}$$

$\rho_o, P_o$



# Nondimensional Ballistic Coefficient

$$\frac{v}{v_e} = \exp\left(\frac{h_s \rho_o}{2\beta \sin \gamma} \frac{\rho}{\rho_o}\right) = \exp\left(\frac{P_o}{2\beta g \sin \gamma} \frac{\rho}{\rho_o}\right)$$

Let  $\hat{\beta} \equiv \frac{\beta g}{P_o}$  (Nondimensional form of ballistic coefficient)

Note that we are using the estimated value of  $P_o = \rho_o g h_s$ , not the actual surface pressure.

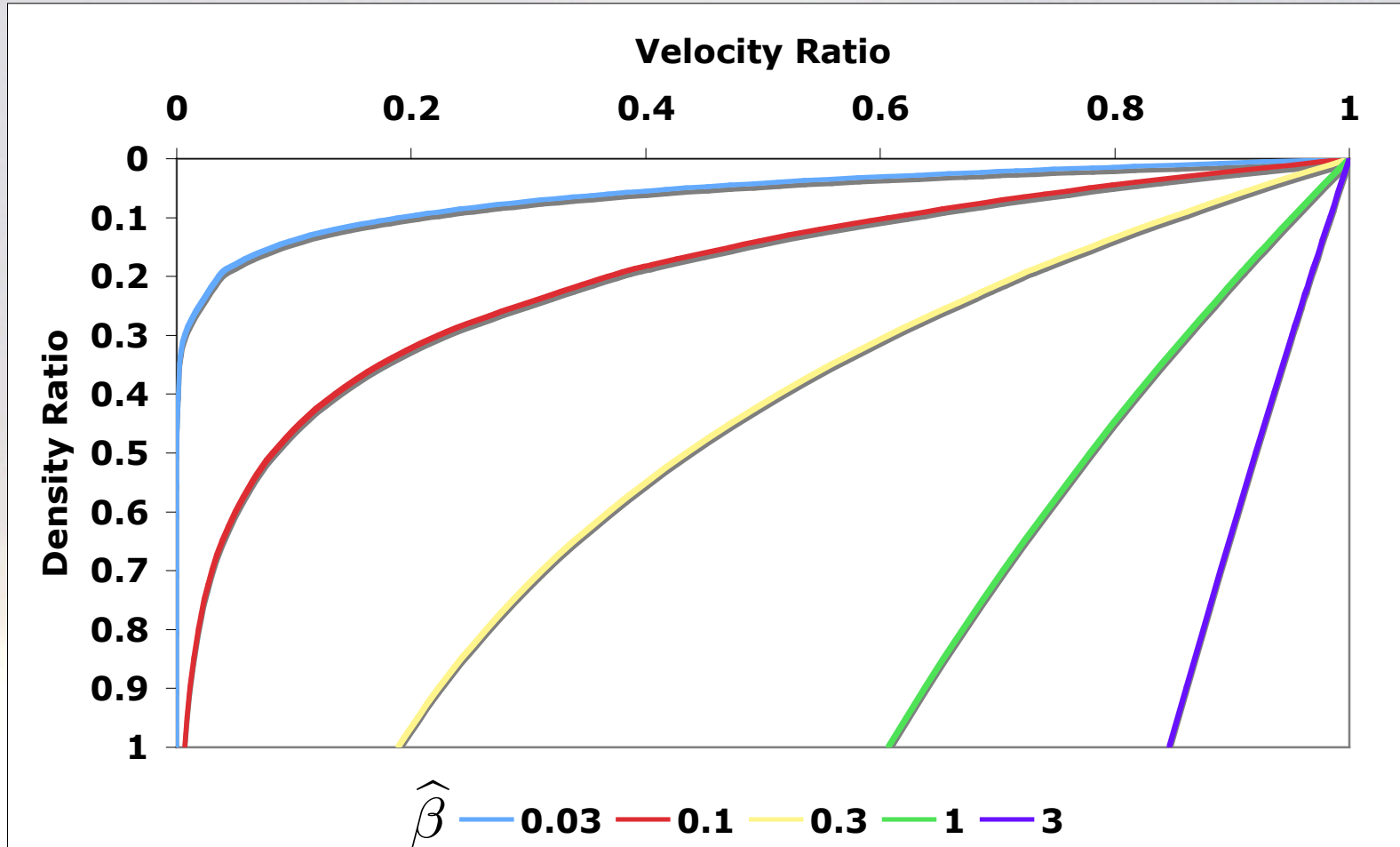
$$\frac{v}{v_e} = \exp\left(\frac{1}{2\hat{\beta} \sin \gamma} \frac{\rho}{\rho_o}\right)$$

$$\beta_{crit} = -\frac{\rho_o h_s}{\sin \gamma}$$

$$\hat{\beta}_{crit} = -\frac{1}{\sin \gamma}$$



# Entry Velocity Trends, $\gamma = -90^\circ$



# Ballistic Entry (no lift)

$s$  = distance along the flight path

$$\frac{dv}{dt} = -g \sin \gamma - \frac{D}{m}$$

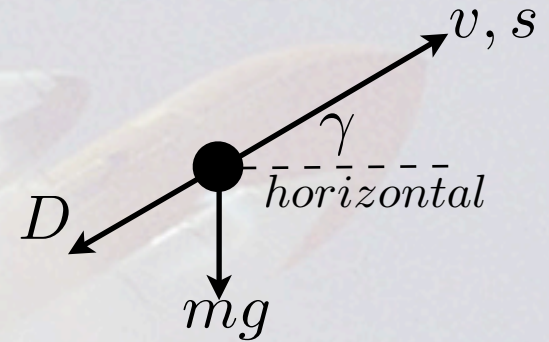
Again assuming  $D \gg g$ ,

$$\frac{dv}{dt} = -\frac{D}{m} \quad \text{Drag } D \equiv \frac{1}{2} \rho v^2 A c_D$$

$$\frac{dv}{dt} = -\frac{\rho c_D A}{2m} v^2$$

Separating the variables,

$$\frac{dv}{v^2} = -\frac{\rho}{2\beta} dt$$



# Calculating the Entry Velocity Profile

$$\frac{dh}{dt} = v \sin \gamma \Rightarrow dt = \frac{dh}{v \sin \gamma}$$

$$\frac{dv}{v^2} = -\frac{\rho}{2\beta v \sin \gamma} dh \Rightarrow \frac{dv}{v} = -\frac{\rho}{2\beta \sin \gamma} dh$$

$$\frac{dv}{v} = -\frac{\rho_o}{2\beta \sin \gamma} e^{-\frac{h}{h_s}} dh$$

$$\int_{v_e}^v \frac{dv}{v} = -\frac{\rho_o}{2\beta \sin \gamma} \int_{h_e}^h e^{-\frac{h}{h_s}} dh$$

$$\ln \frac{v}{v_e} = \frac{\rho_o h_s}{2\beta \sin \gamma} \left[ e^{-\frac{h}{h_s}} \right]_{h_e}^h = \frac{1}{2\hat{\beta} \sin \gamma} \left[ e^{-\frac{h}{h_s}} - e^{-\frac{h_e}{h_s}} \right]$$



# Deriving the Entry Velocity Function

Remember that  $e^{-\frac{h_e}{h_s}} = \frac{\rho_e}{\rho_o} \approx 0$

$$\frac{v}{v_e} = \exp \left[ \frac{1}{2\hat{\beta} \sin \gamma} e^{-\frac{h}{h_s}} \right]$$

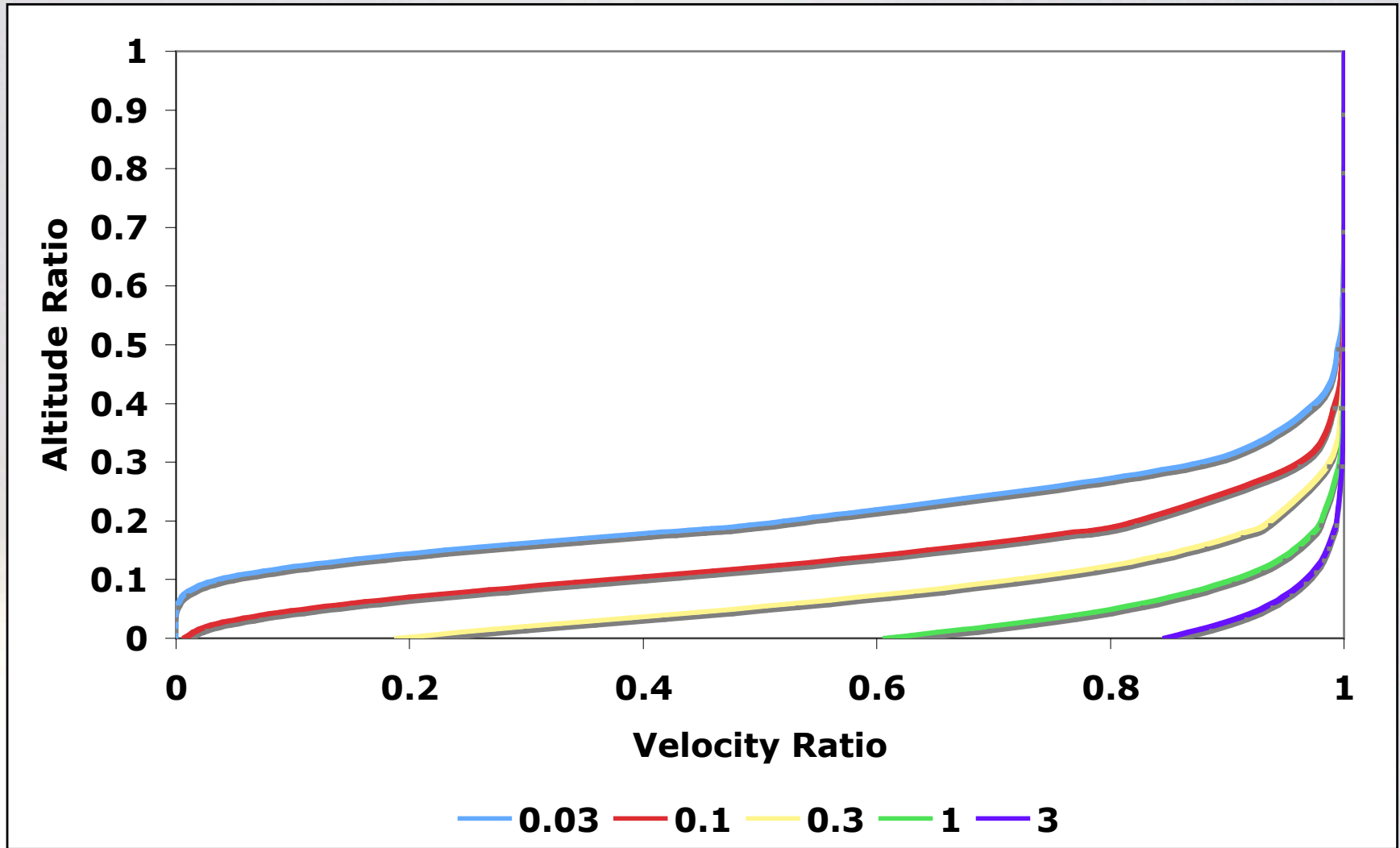
We have a parametric entry equation in terms of nondimensional velocity ratios, ballistic coefficient, and altitude. To bound the nondimensional altitude variable between 0 and 1, rewrite as

$$\frac{v}{v_e} = \exp \left[ \frac{1}{2\hat{\beta} \sin \gamma} e^{-\frac{h}{h_e} \frac{h_e}{h_s}} \right]$$

$\frac{h_e}{h_s}$  and  $\hat{\beta}$  are the only variables that relate to a specific planet



# Earth Entry, $\gamma = -90^\circ$



# Deceleration as a Function of Altitude

Start with

$$\frac{v}{v_e} = \exp \left[ \frac{1}{2\hat{\beta} \sin \gamma} e^{-\frac{h}{h_e} \frac{h_e}{h_s}} \right] \quad \text{Let } B \equiv \frac{1}{2\hat{\beta} \sin \gamma}$$

$$\frac{v}{v_e} = \exp \left( B e^{-\frac{h}{h_s}} \right)$$

$$\frac{d}{dt} \left( \frac{v}{v_e} \right) = \exp \left( B e^{-\frac{h}{h_s}} \right) \frac{d}{dt} \left( B e^{-\frac{h}{h_s}} \right)$$

$$\frac{dv}{dt} = v_e \exp \left( B e^{-\frac{h}{h_s}} \right) \frac{-B}{h_s} \left( e^{-\frac{h}{h_s}} \right) \frac{dh}{dt}$$

$$\frac{dh}{dt} = v \sin \gamma = v_e \sin \gamma \exp B e^{-\frac{h}{h_s}}$$



# Parametric Deceleration

$$\frac{dv}{dt} = v_e \exp\left(Be^{-\frac{h}{h_s}}\right) \frac{-B}{h_s} \left(e^{-\frac{h}{h_s}}\right) v_e \sin \gamma \exp\left(Be^{-\frac{h}{h_s}}\right)$$

$$\frac{dv}{dt} = \frac{-Bv_e^2}{h_s} \sin \gamma \left(e^{-\frac{h}{h_s}}\right) \exp\left(2Be^{-\frac{h}{h_s}}\right)$$

$$\frac{dv}{dt} = \frac{-v_e^2}{2h_s \hat{\beta}} \left(e^{-\frac{h}{h_s}}\right) \exp\left(\frac{1}{\hat{\beta} \sin \gamma} e^{-\frac{h}{h_s}}\right)$$

$$\text{Let } n_{ref} \equiv \frac{v_e^2}{h_s}, \nu \equiv \frac{dv/dt}{n_{ref}}, \varphi \equiv \frac{h_e}{h_s}$$

$$\nu = \frac{-1}{2\hat{\beta}} \left(e^{-\varphi \frac{h}{h_e}}\right) \exp\left(\frac{1}{\hat{\beta} \sin \gamma} e^{-\varphi \frac{h}{h_e}}\right)$$



# Nondimensional Deceleration, $\gamma = -90^\circ$

