

# Entry Aerodynamics

- Atmospheric Regimes on Entry
- Basic fluid parameters
- Definition of Mean Free Path
- Rarified gas Newtonian flow
- Continuum Newtonian flow (hypersonics)



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# Basic Fluids Parameters

$$M \equiv \text{Mach Number} = \frac{v}{a}$$

$$a \equiv \text{speed of sound} = \sqrt{\gamma RT} \quad \left( R = \frac{\mathfrak{R}}{\bar{m}} \right)$$

$$\frac{\text{ordered energy}}{\text{random energy}} = \frac{\frac{1}{2}mv^2}{\frac{1}{2}m\bar{v}_g^2} = \frac{v^2}{3RT} = \frac{\gamma v^2}{3a^2} = \frac{\gamma}{3}M^2$$

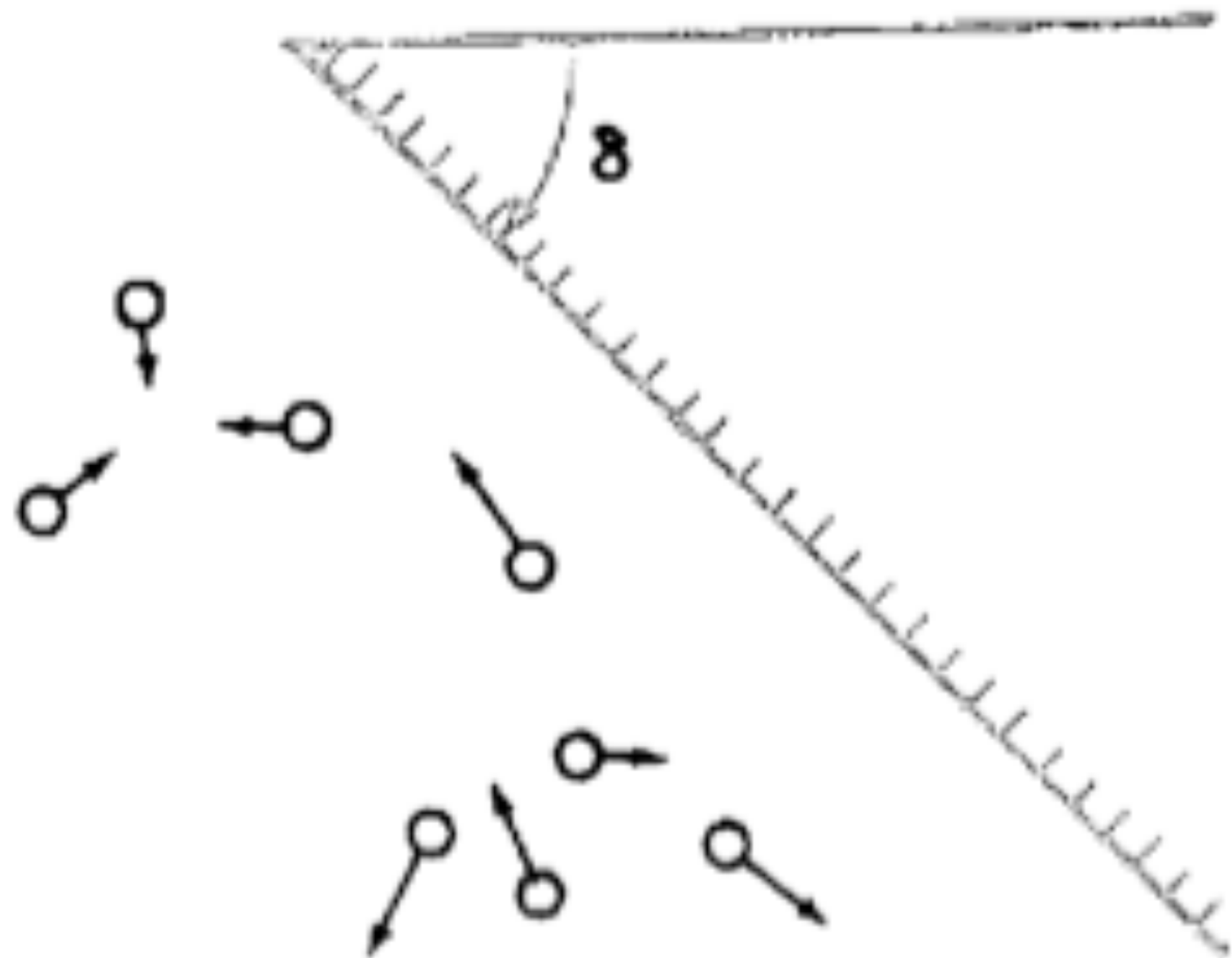
$$Re \equiv \text{Reynold's number} = \frac{\text{inertial force}}{\text{viscous force}}$$

$$Re = \frac{\dot{m}v}{\tau A} = \frac{\rho Av^2}{\mu \frac{v}{L} A} = \frac{\rho v L}{\mu}$$

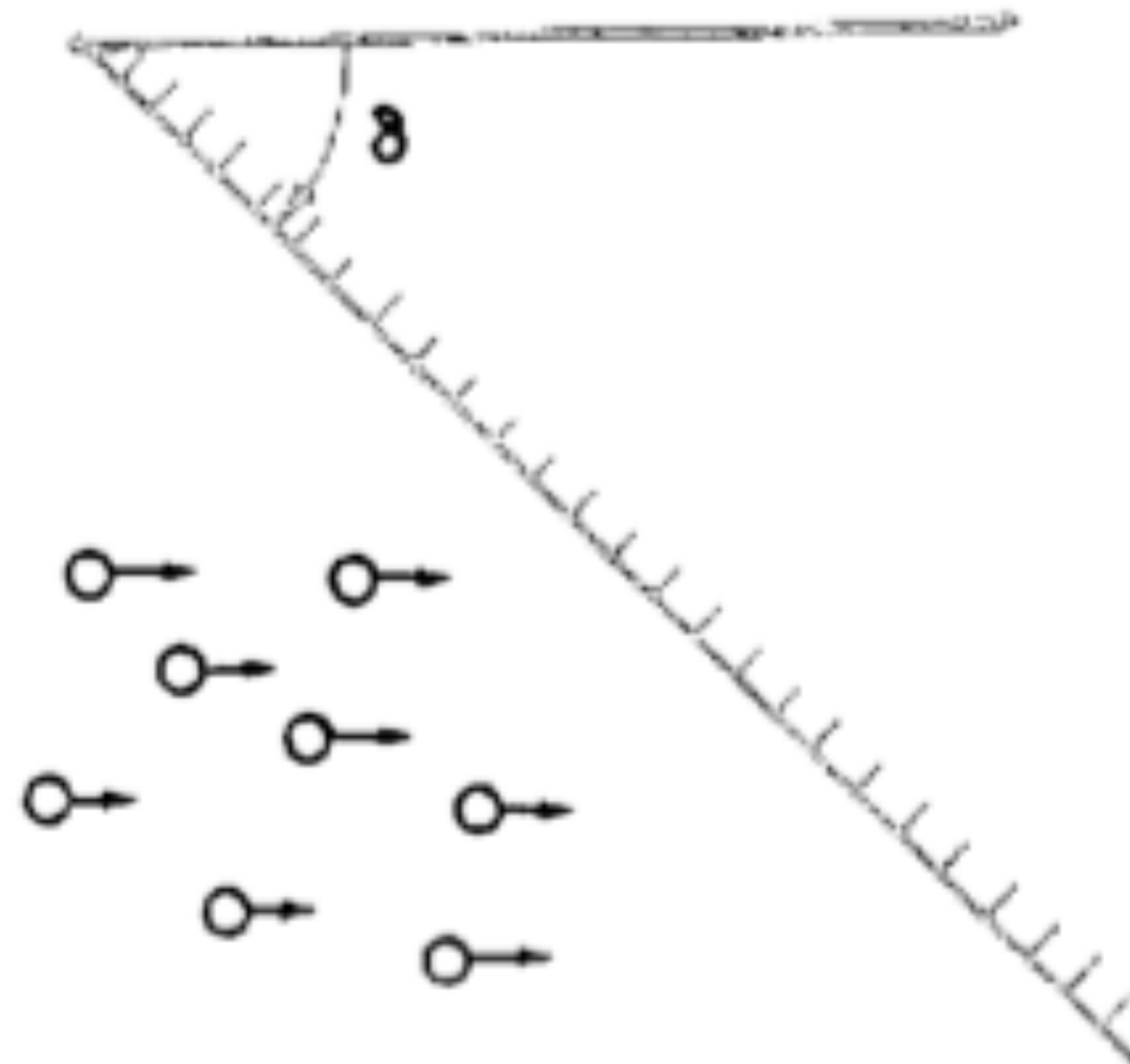




# Random vs. Ordered Energy



$M < 1$



$M > 1$





# More Fluid Parameters

$K \equiv$  Knudsen number

$$K = \frac{\text{number of collisions with body}}{\text{number of collisions with other molecules}}$$

$$K = \frac{\lambda}{L}$$

$\lambda \equiv$  mean free path

$L \equiv$  vehicle characteristic length





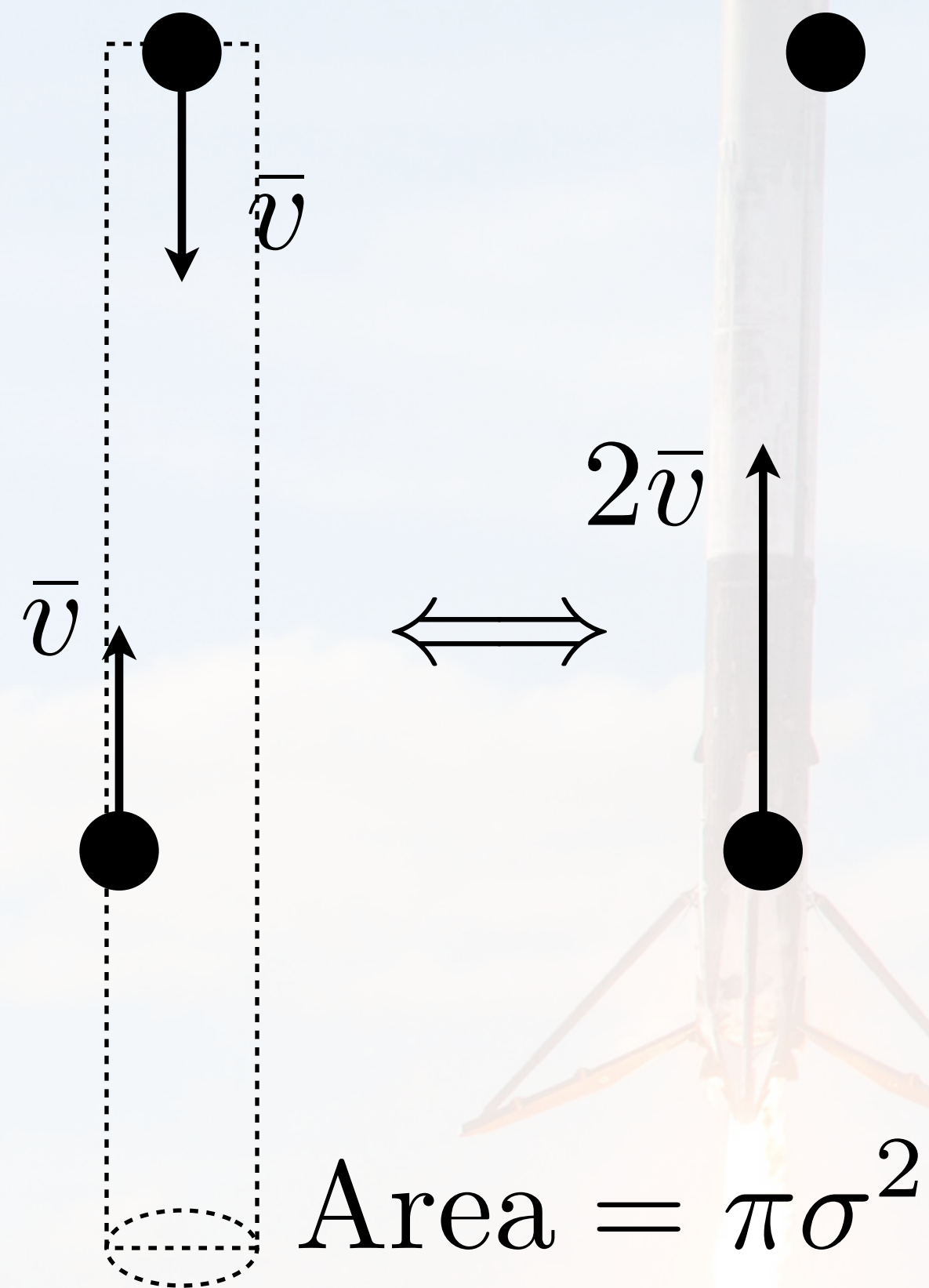
# Estimating Mean Free Path

Assume:

- All molecules are perfect rigid spheres
- Each has diameter  $\sigma$ , mass  $m$ , and velocity  $\bar{v}$
- Consider a cube with side length  $L$  containing  $N$  molecules
- $N/6$  molecules are traveling in each direction
  - $\pm X$
  - $\pm Y$
  - $\pm Z$



# Consider Collisions in +Z Direction



- number of potential +Z collisions

$$n(+Z) = \frac{1}{6} N \frac{\pi\sigma^2 L}{L^3} = \frac{1}{6} N \frac{\pi\sigma^2}{L^2}$$

frequency of +Z collisions

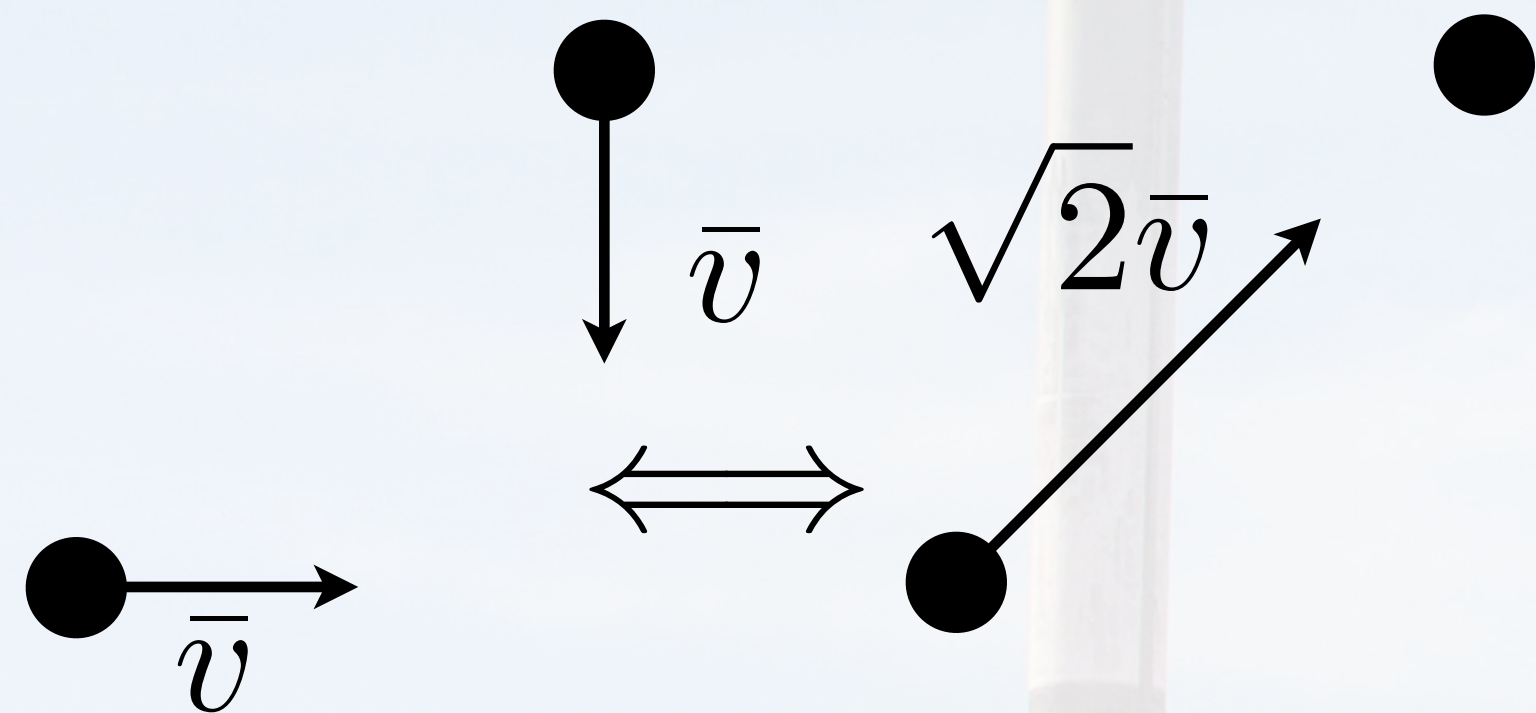
$$f(+Z) = \frac{n(+Z)}{\Delta t} = \frac{\frac{\pi}{6} N \frac{\sigma^2}{L^2}}{\frac{L}{2\bar{v}}}$$

$$f(+Z) = \frac{\pi\rho\sigma^2\bar{v}}{3m}$$





# Consider Collisions in +X Direction



frequency of +X collisions

$$f(+X) = \frac{n(+X)}{\Delta t} = \frac{\sqrt{2}\pi\rho\sigma^2\bar{v}}{6m}$$

$$f(-X) = f(+Y) = f(-Y) = f(+X) \quad f(-Z) = 0$$

Total frequency of collisions

$$f = \frac{\pi}{3}(1 + 2\sqrt{2})\frac{\rho\sigma^2\bar{v}}{m}$$



# Mean Free Path

$$\lambda = \frac{\bar{v}}{f} = \frac{m/\sigma^2}{\frac{\pi}{3}(1+2\sqrt{2})\rho} \quad \left( \propto \frac{1}{\rho} \right)$$

at sea level:  $\lambda = 6.7 \times 10^{-8} \text{ m}$

at 100 km:  $\lambda = 0.3 \text{ m} \sim 1 \text{ ft}$





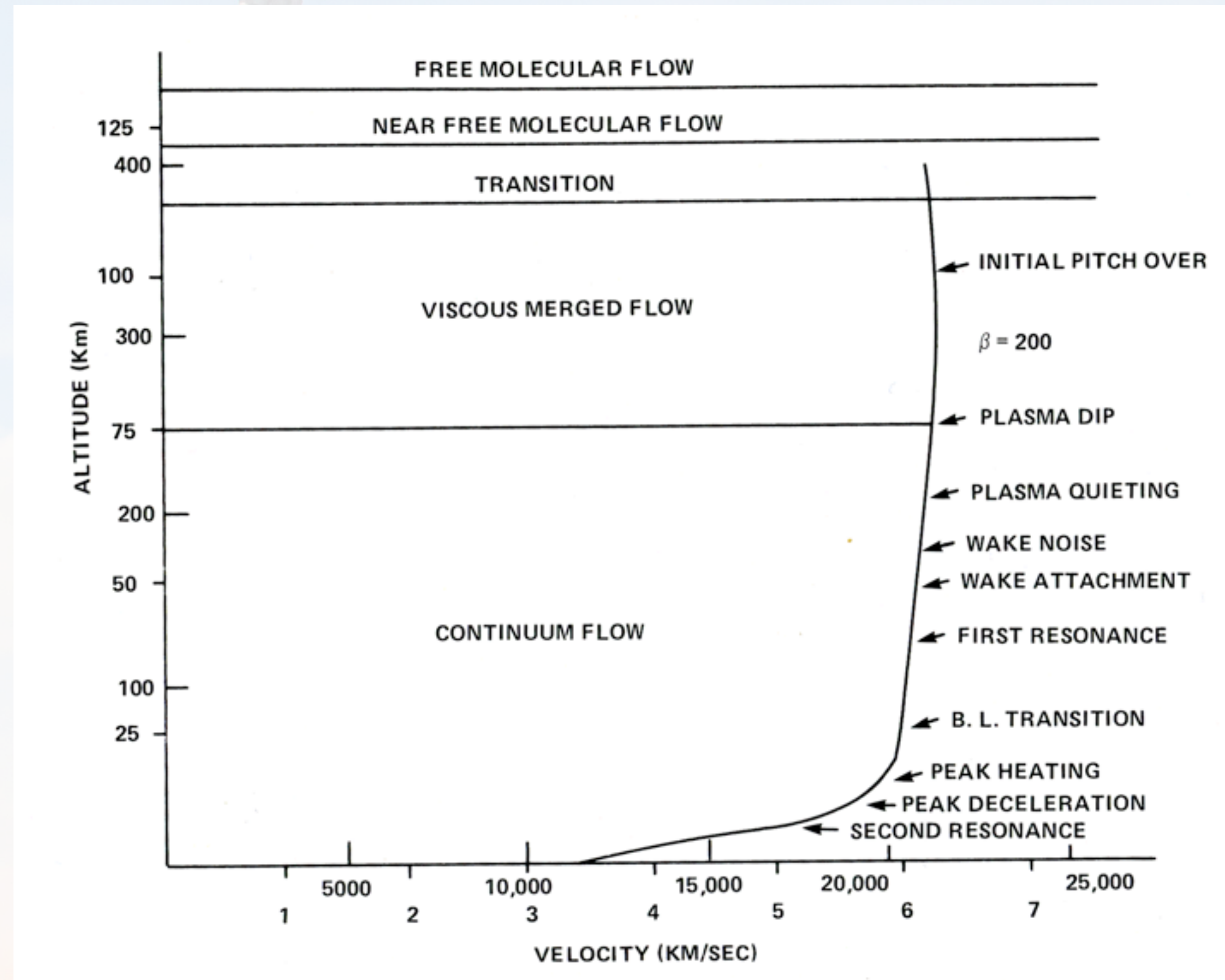
# Five Basic Flow Regimes

- Free molecular regime
- Near-free molecular regime
- Transition regime
- Viscous merged boundary layer
- Continuous regime





# Entry Flow Regimes



ref: Frank J. Regan, Reentry Vehicle Dynamics AIAA Education Series, NY, NY 1984





# Flow Regime Definitions

Knudsen number in rarified flow

$$K = \frac{\lambda}{R_N}$$

Mean free path after collision

$$\lambda_c = \frac{4}{\sqrt{\pi\gamma}} \left( \frac{T_w}{T_\infty} \right)^{\frac{1}{2}} \frac{\lambda_\infty}{M_\infty}$$

If  $T_w \sim T_\infty$

$$\lambda_c \cong 1.9 \frac{\lambda_\infty}{M_\infty}$$





# Free Molecular Regime

- Orbital flight

$$\lambda \gg \ell$$

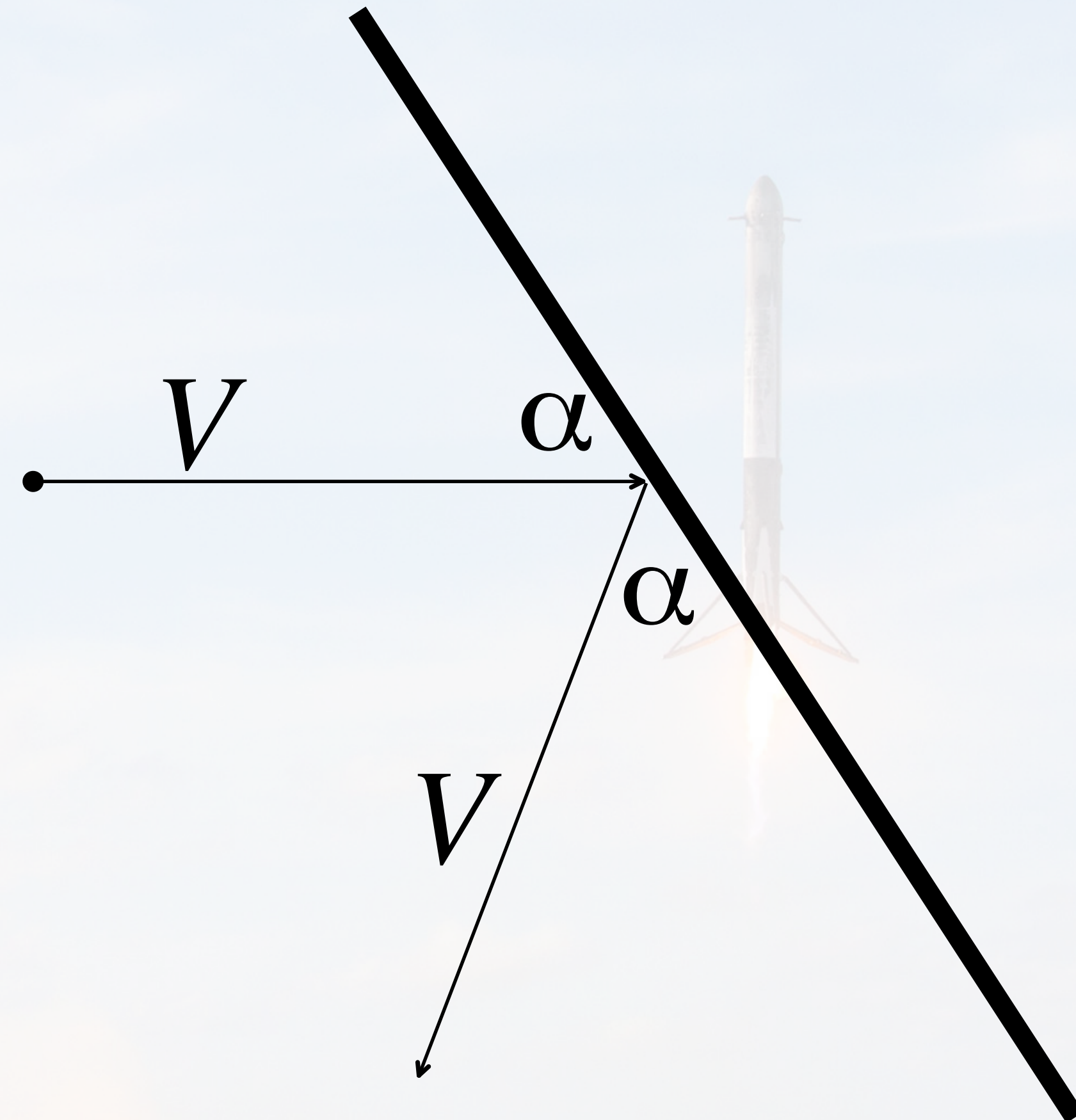
- Molecule encountering a boundary (e.g., surface of vehicle) attains the state of the boundary after a single collision

$$K_c \geq 10 \text{ or } K_\infty > 5.24M_\infty$$



# Newtonian Flow

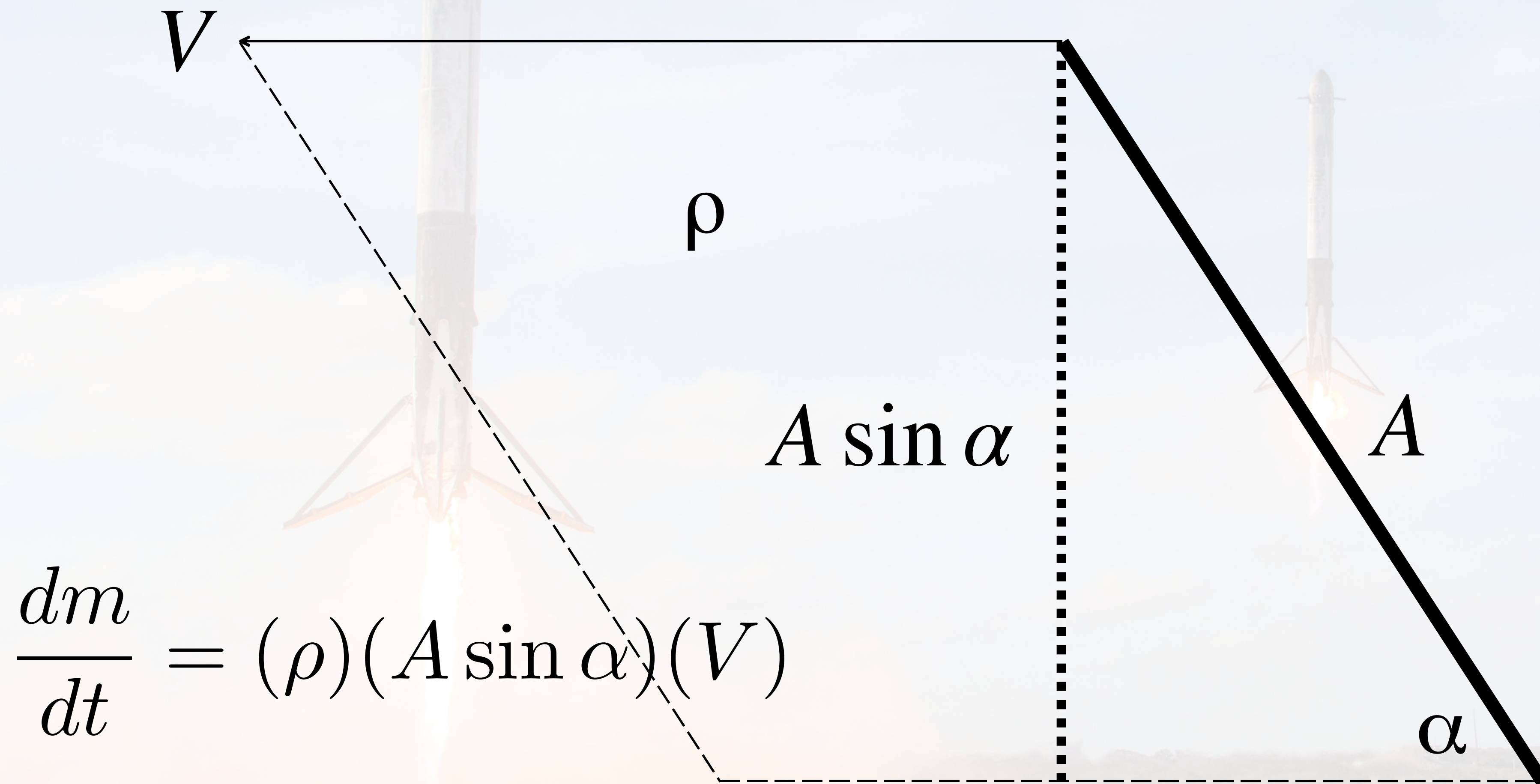
- Mean free path of particles much larger than spacecraft --> no appreciable interaction of air molecules
- Model vehicle / atmosphere interactions as independent perfectly elastic collisions





# Newtonian Analysis

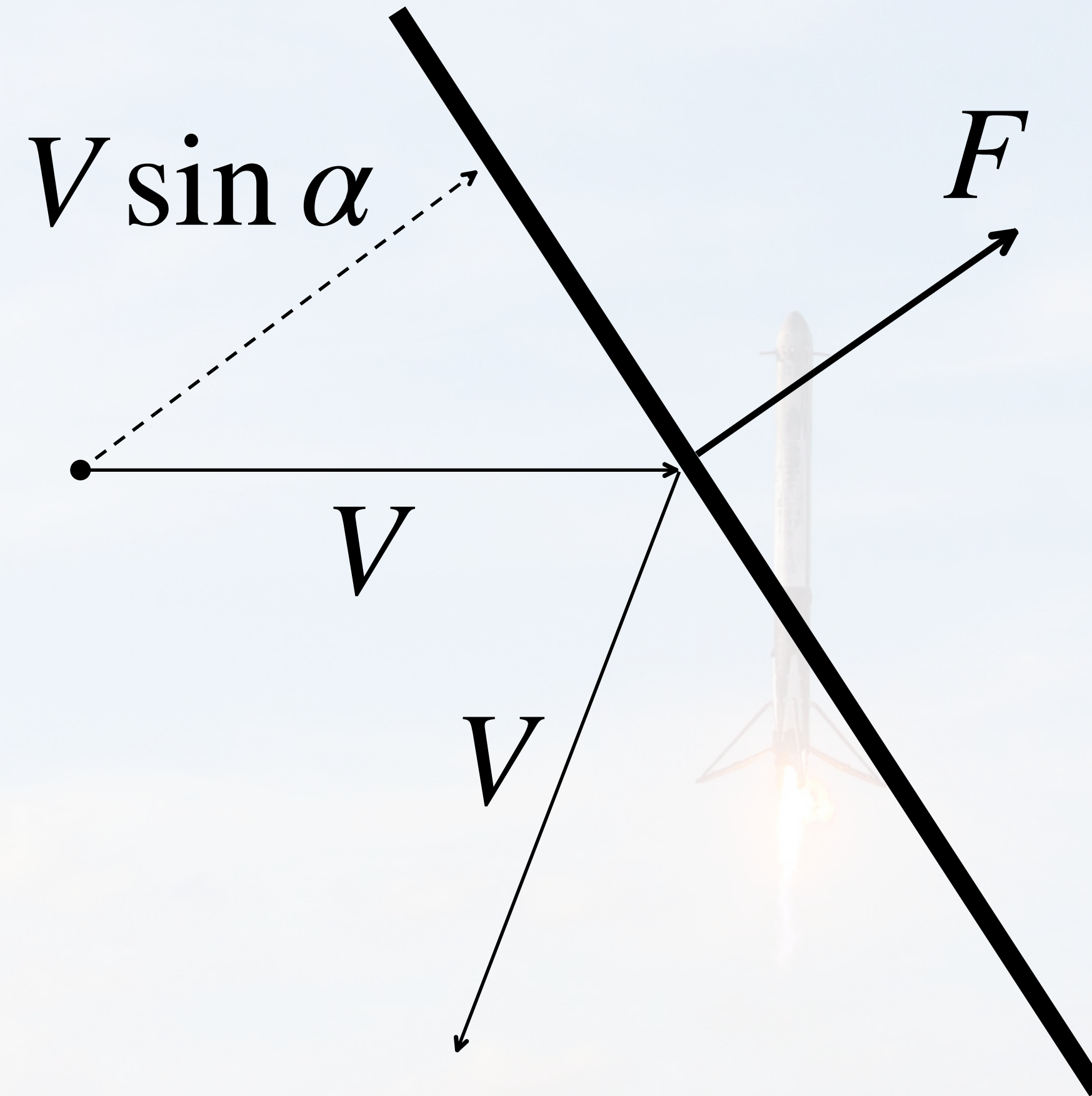
mass flux = (density)(swept area)(velocity)





# Momentum Transfer

- Momentum perpendicular to wall is reversed at impact
- “Bounce” momentum is transferred to vehicle
- Momentum parallel to wall is unchanged



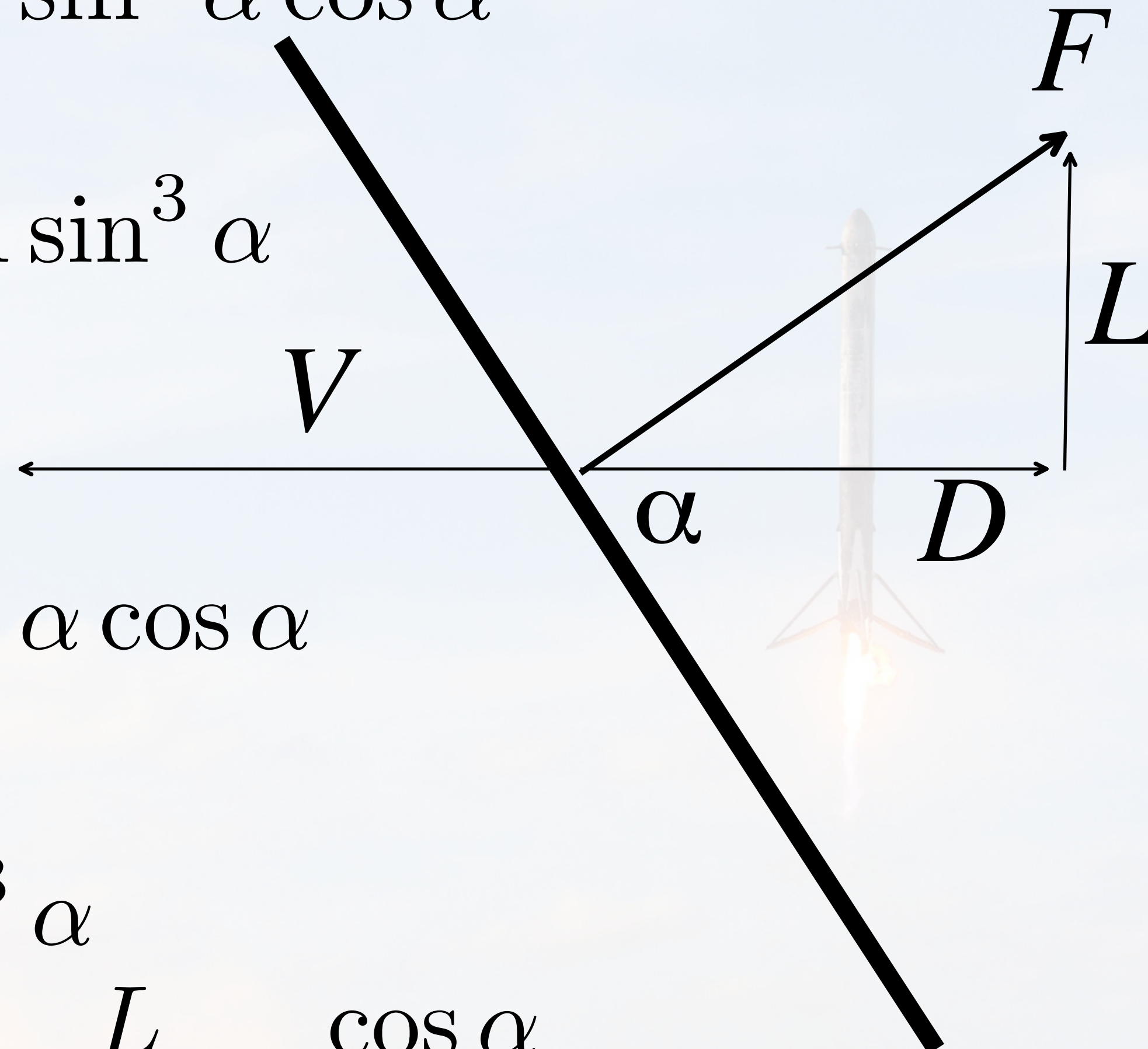
$$F = \frac{dm}{dt} \Delta V = \rho V A \sin \alpha (2V \sin \alpha) = 2\rho V^2 A \sin^2 \alpha$$



# Lift and Drag

$$L = F \cos \alpha = 2\rho V^2 A \sin^2 \alpha \cos \alpha$$

$$D = F \sin \alpha = 2\rho V^2 A \sin^3 \alpha$$



$$c_L = \frac{L}{\frac{1}{2}\rho V^2 A} = 4 \sin^2 \alpha \cos \alpha$$

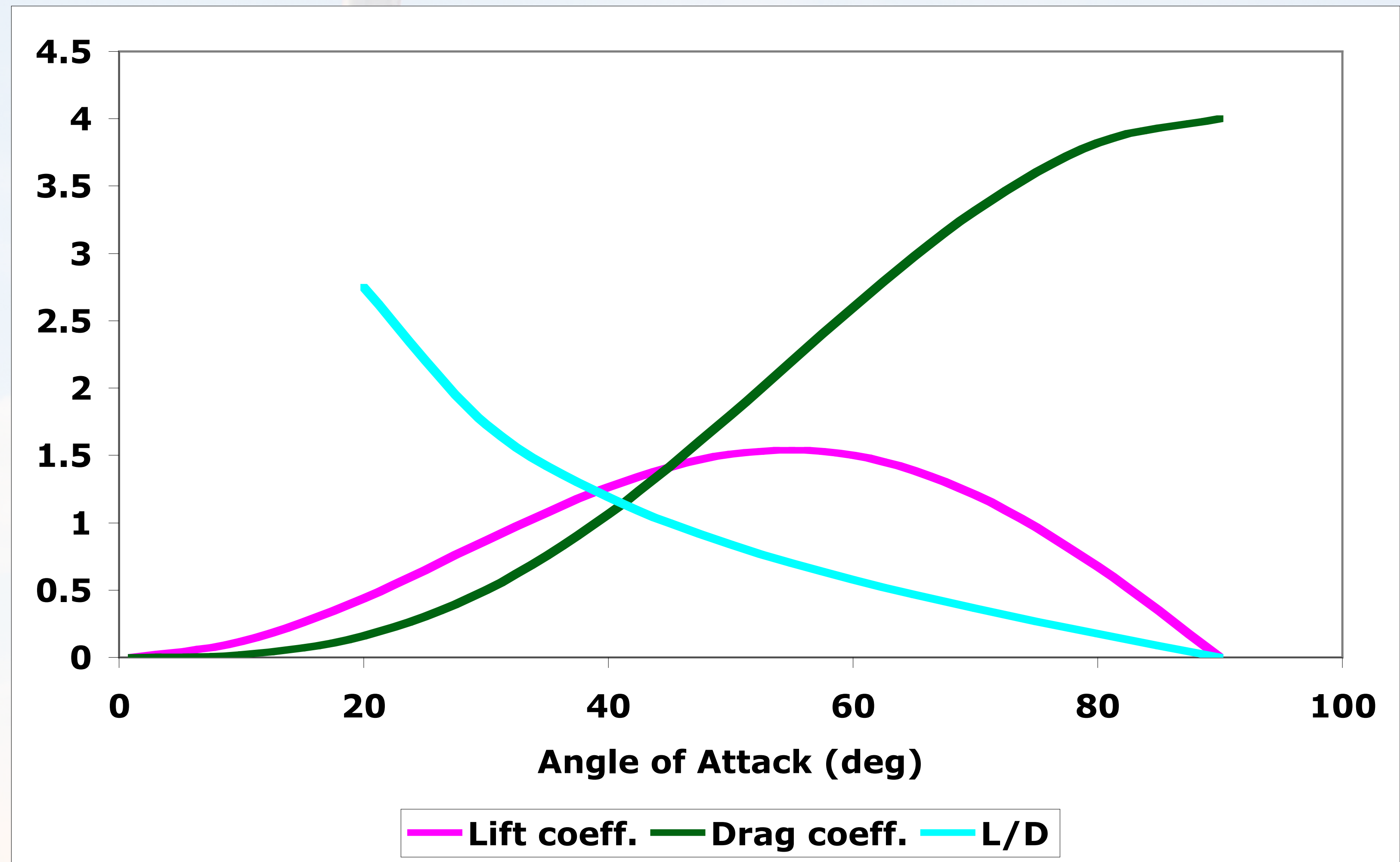
$$c_D = \frac{D}{\frac{1}{2}\rho V^2 A} = 4 \sin^3 \alpha$$

$$\frac{L}{D} = \frac{\cos \alpha}{\sin \alpha} = \cot \alpha$$





# Flat Plate Newtonian Aerodynamics





# Example of Newtonian Flow Calculations

Consider a cylinder of length  $l$ , entering atmosphere transverse to flow

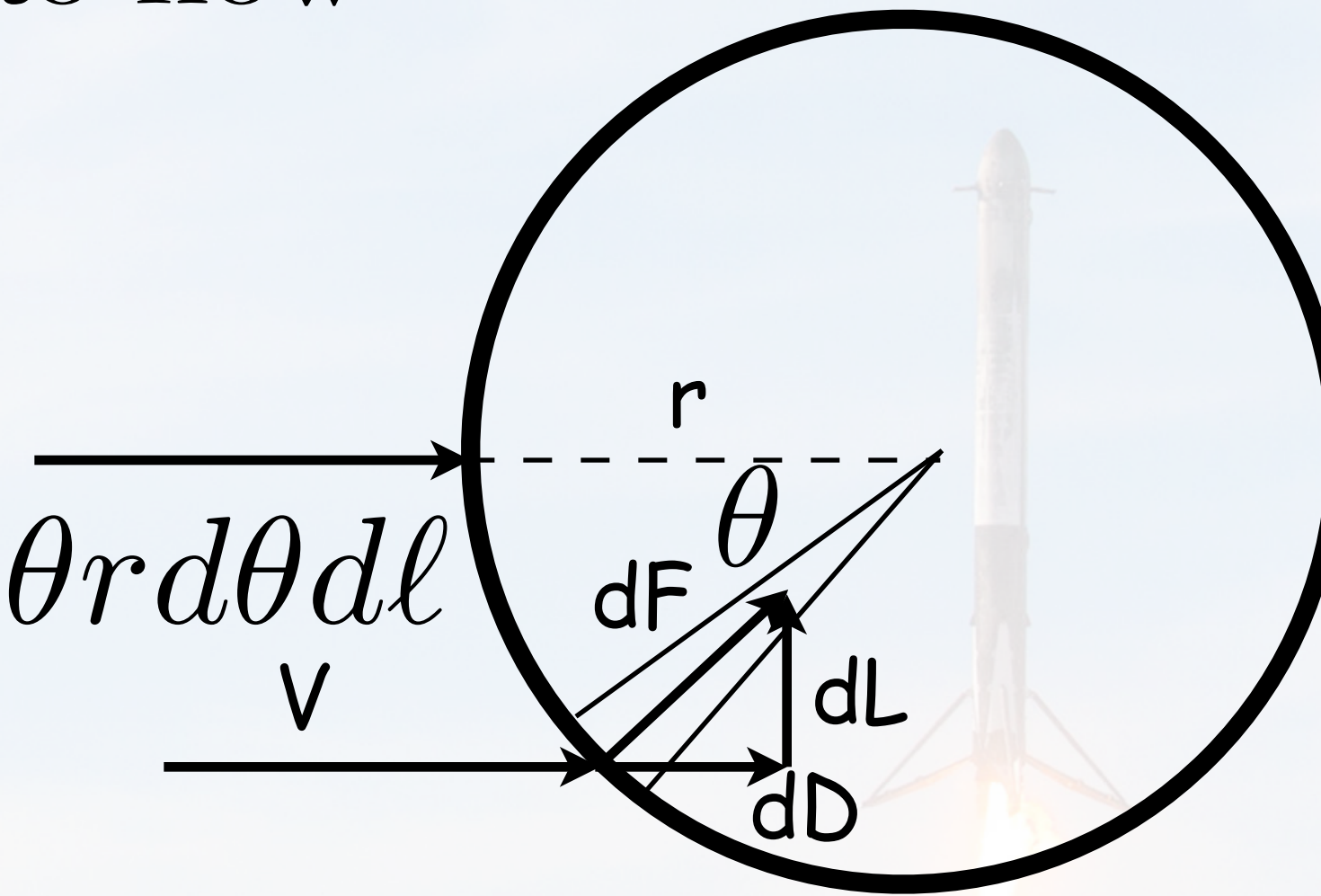
$$dA = r d\theta dl$$

$$d\dot{m} = \rho dA \cos \theta V = \rho V \cos \theta r d\theta dl$$

$$dF = d\dot{m} \Delta V = 2\rho V^2 \cos^2 \theta r d\theta dl$$

$$dD = dF \cos \theta = 2\rho V^2 \cos^3 \theta r d\theta dl$$

$$dL = dF \sin \theta = 2\rho V^2 \cos^2 \theta \sin \theta r d\theta dl$$





# Integration to Find Drag Coefficient

Integrate from  $\theta = -\frac{\pi}{2} \rightarrow \frac{\pi}{2}$

$$D = \int_{-\frac{\pi}{2}}^{+\frac{\pi}{2}} \int_0^{\ell} dD = 2\rho V^2 r \int_{-\frac{\pi}{2}}^{+\frac{\pi}{2}} \int_0^{\ell} \cos^3 \theta d\theta d\ell$$

$$= 2\rho V^2 r \ell \int_{-\frac{\pi}{2}}^{+\frac{\pi}{2}} \cos^3 \theta d\theta = \frac{8}{3} \rho V^2 r \ell$$

By definition,  $D = \frac{1}{2} \rho V^2 A c_D$  and, for a cylinder  $A = 2r\ell$

$$\rho V^2 r \ell c_D = \frac{8}{3} \rho V^2 r \ell \implies c_D = \frac{8}{3}$$





# Near Free Molecular Flow Regime

- Also known as “slip region”
- Gas molecule only attains state of moving boundary after several collisions
- Molecules near the wall will have a different velocity from the wall
- Temperature will be nearly discontinuous function of separation from wall

$$10 \leq K_c \leq \frac{1}{3} \text{ or } 5.4M_\infty \leq K_\infty \leq 0.175M_\infty$$



# Transition Region

- Very difficult to treat analytically
- For engineering purposes, usually treated as interpolation between slip and viscous flow

$$0.175M_{\infty} \leq K_{\infty} \leq 1$$



# Viscous Merged Layer Regime

- Viscous effects in forming shock and boundary layer must be treated in a unified manner
  - Boundary layer on the wall alters the conditions for the forming shock wave
  - Large pressure gradients across the shock wave significantly alter the boundary layer
- Neither shocks nor boundary layers can be treated as discontinuities

$$1 \leq K_{\infty} \leq \frac{0.1}{\rho_s / \rho_{\infty}}$$



# Continuous Regime

- Classical fluid mechanics of high Reynolds number
- Shock waves and boundary layer treated as discontinuities

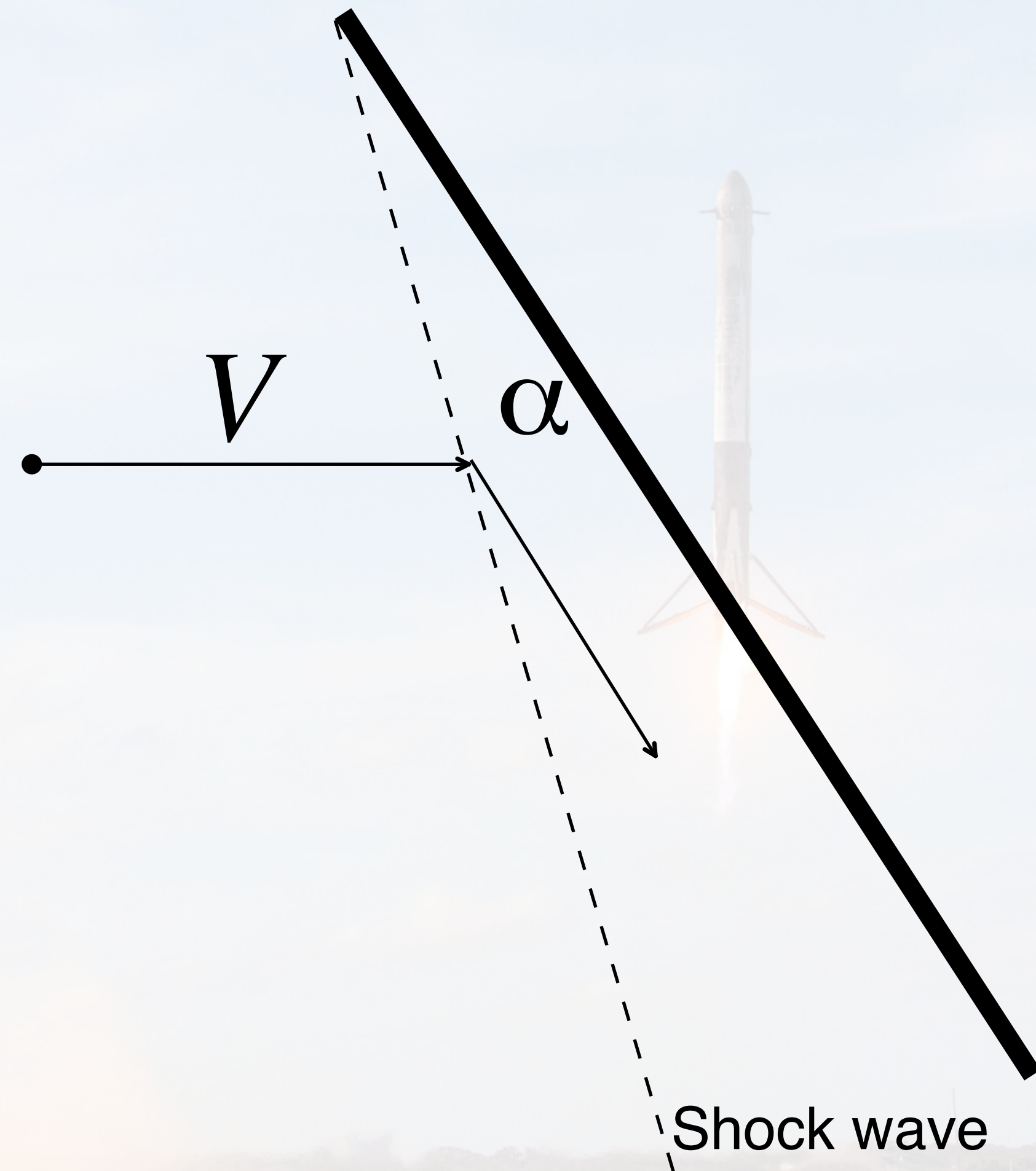
$$K_{\infty} > \frac{0.1}{\rho_s / \rho_{\infty}}$$

- Subdivided based on Mach number
  - Incompressible (subsonic) ( $M \leq \sim 0.8$ )
  - Transonic ( $\sim 0.8 \leq M \leq \sim 1.3$ )
  - Supersonic ( $\sim 1.3 \leq M \leq \sim 5$ )
  - Hypersonic ( $\sim 5 \leq M$ )



# Continuum Newtonian Flow (Hypersonics)

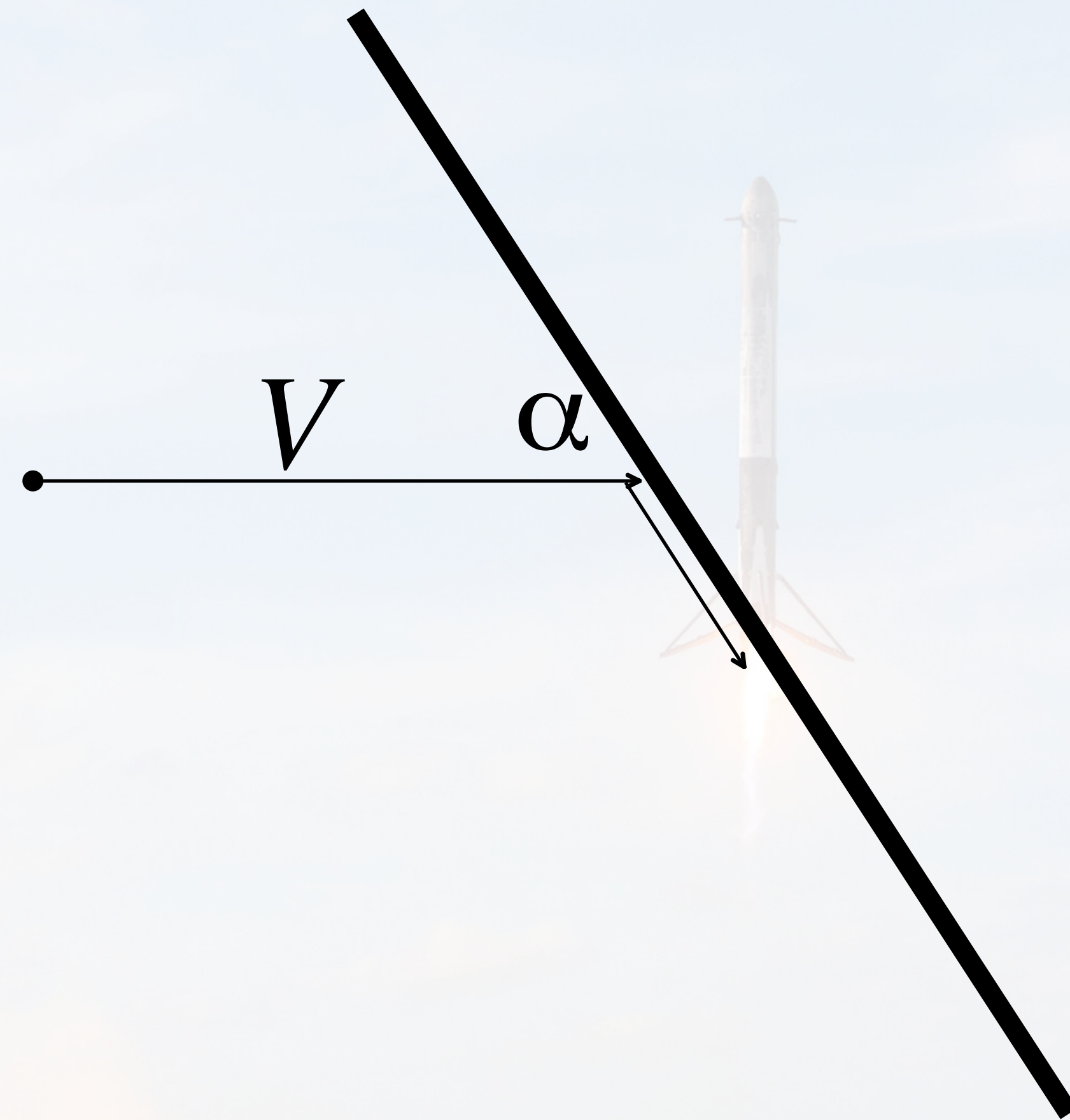
- Air molecules predominately interact with shock waves
- Effect of shock wave passage is to decelerate flow and turn it parallel to vehicle surface





# Continuum Newtonian Flow (Hypersonics)

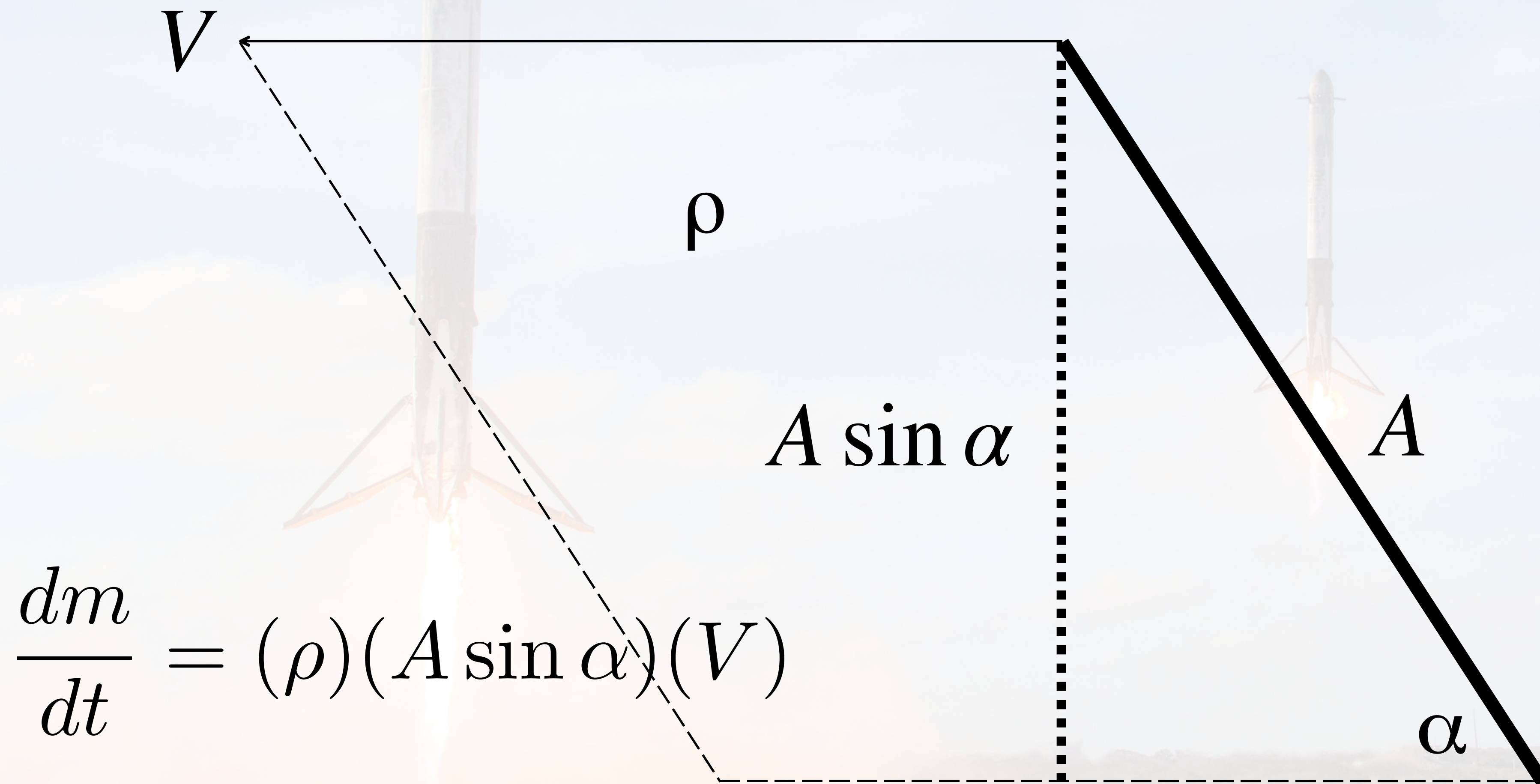
- Treat hypersonic aerodynamics in manner similar to previous Newtonian flow analysis
- All momentum perpendicular to wall is absorbed by the wall





# Mass Flux (unchanged)

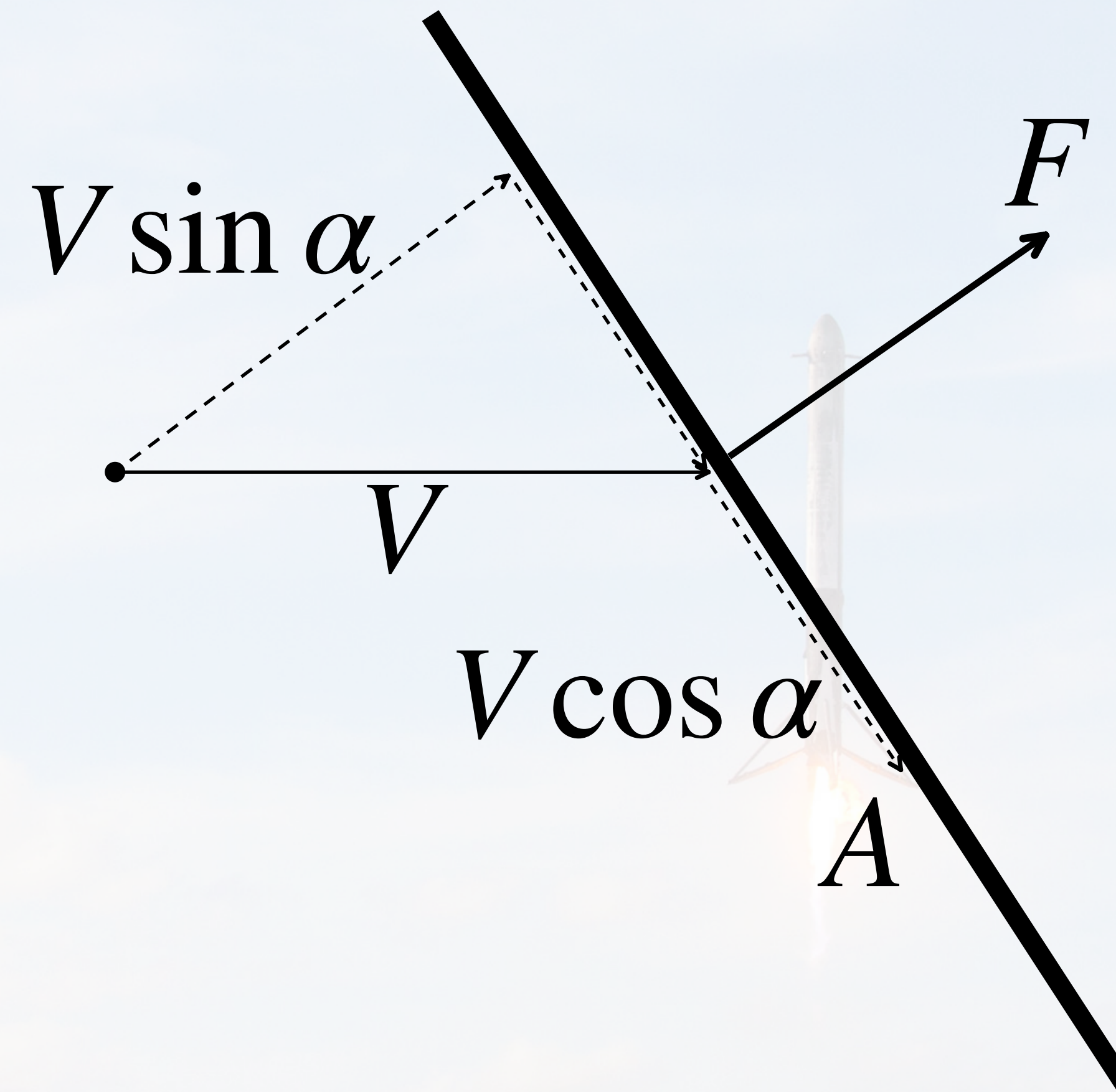
mass flux = (density)(swept area)(velocity)





# Momentum Transfer

- Momentum perpendicular to wall is absorbed at impact and transferred to vehicle
- Momentum parallel to wall is unchanged



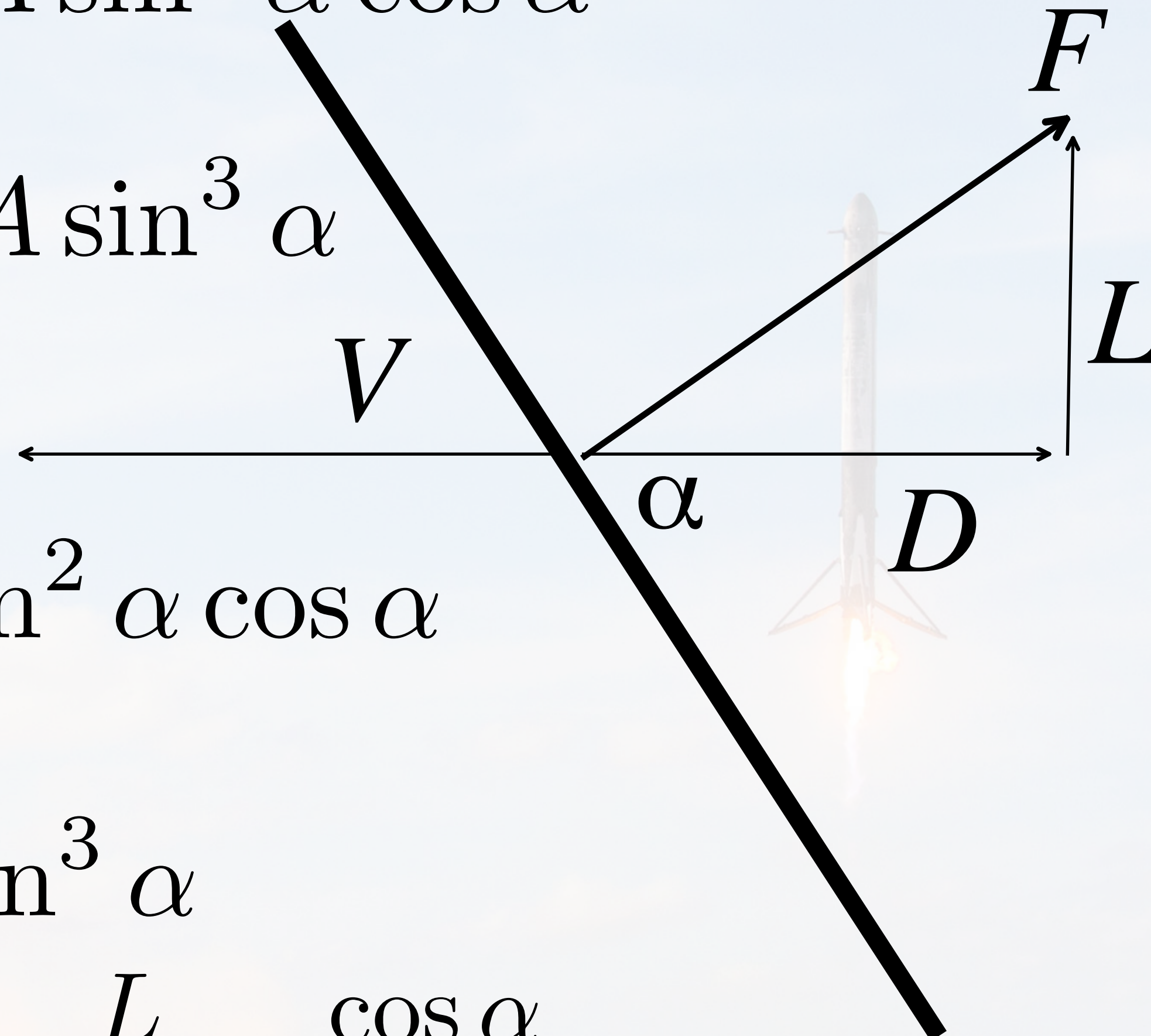
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# Lift and Drag

$$L = F \cos \alpha = \rho V^2 A \sin^2 \alpha \cos \alpha$$

$$D = F \sin \alpha = \rho V^2 A \sin^3 \alpha$$



$$c_L = \frac{L}{\frac{1}{2} \rho V^2 A} = 2 \sin^2 \alpha \cos \alpha$$

$$c_D = \frac{D}{\frac{1}{2} \rho V^2 A} = 2 \sin^3 \alpha$$

$$\frac{L}{D} = \frac{\cos \alpha}{\sin \alpha} = \cot \alpha$$





# Modified Newtonian Flow

- Coefficient of pressure in “classical” Newtonian flow

$$c_p = 2 \sin^2 (\alpha)$$

- Coefficient of pressure in modified Newtonian flow

$$c_p = c_{p_{max}} \sin^2 (\alpha)$$

- $C_p(\max)$  is the pressure coefficient behind a normal shock at flight conditions

$$c_{p_{max}} = \frac{P_{shock} - P_{\infty}}{\frac{1}{2} \rho_{\infty} v_{\infty}^2}$$



# Maximum Coefficient of Pressure

$$C_{p_{max}} = \frac{2}{\gamma M_{\infty}^2} \left\{ \left[ \frac{(\gamma + 1)^2 M_{\infty}^2}{4\gamma M_{\infty}^2 - 2(\gamma - 1)} \right]^{\frac{\gamma}{\gamma-1}} \left[ \frac{1 - \gamma + 2\gamma M_{\infty}^2}{\gamma + 1} \right] - 1 \right\}$$

as  $M \longrightarrow \infty$

$$C_{p_{max}} \longrightarrow \left[ \frac{(\gamma + 1)^2}{4\gamma} \right]^{\frac{\gamma}{\gamma-1}} \left[ \frac{4}{\gamma + 1} \right]$$

$$C_{p_{max}} \longrightarrow 1.839 \text{ for } \gamma = 1.4$$

$$C_{p_{max}} \longrightarrow 2 \text{ for } \gamma = 1$$

