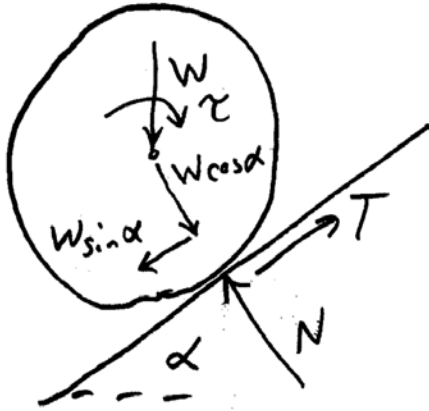


Multi-Wheel Systems

- Obstacle climbing with multiwheel systems
- Planar rocker analysis
- Planar rocker-bogey analysis
- Suspension dynamics



Slopes + Obstacles



Wheel coefficient of friction
with ground = μ

$N \equiv$ normal force to surface

$$\tau = \mu r N = \mu r W \sin \alpha$$

$$T = \mu N = \frac{\tau}{r} = W \sin \alpha$$

$$\mu W \cos \alpha = W \sin \alpha$$

$$\tan \alpha = \mu$$

Assume $\tau > \mu_{\text{limit}} N r$ (friction limited,
not torque limited)

Longitudinal Dynamic Solutions

$$N_1 = mg \left[\left(1 - \frac{a}{\ell} \right) \cos \theta - \left(\frac{h}{\ell} + \frac{r}{\ell} \right) \sin \theta - \frac{1}{g} \frac{dv}{dt} \right]$$

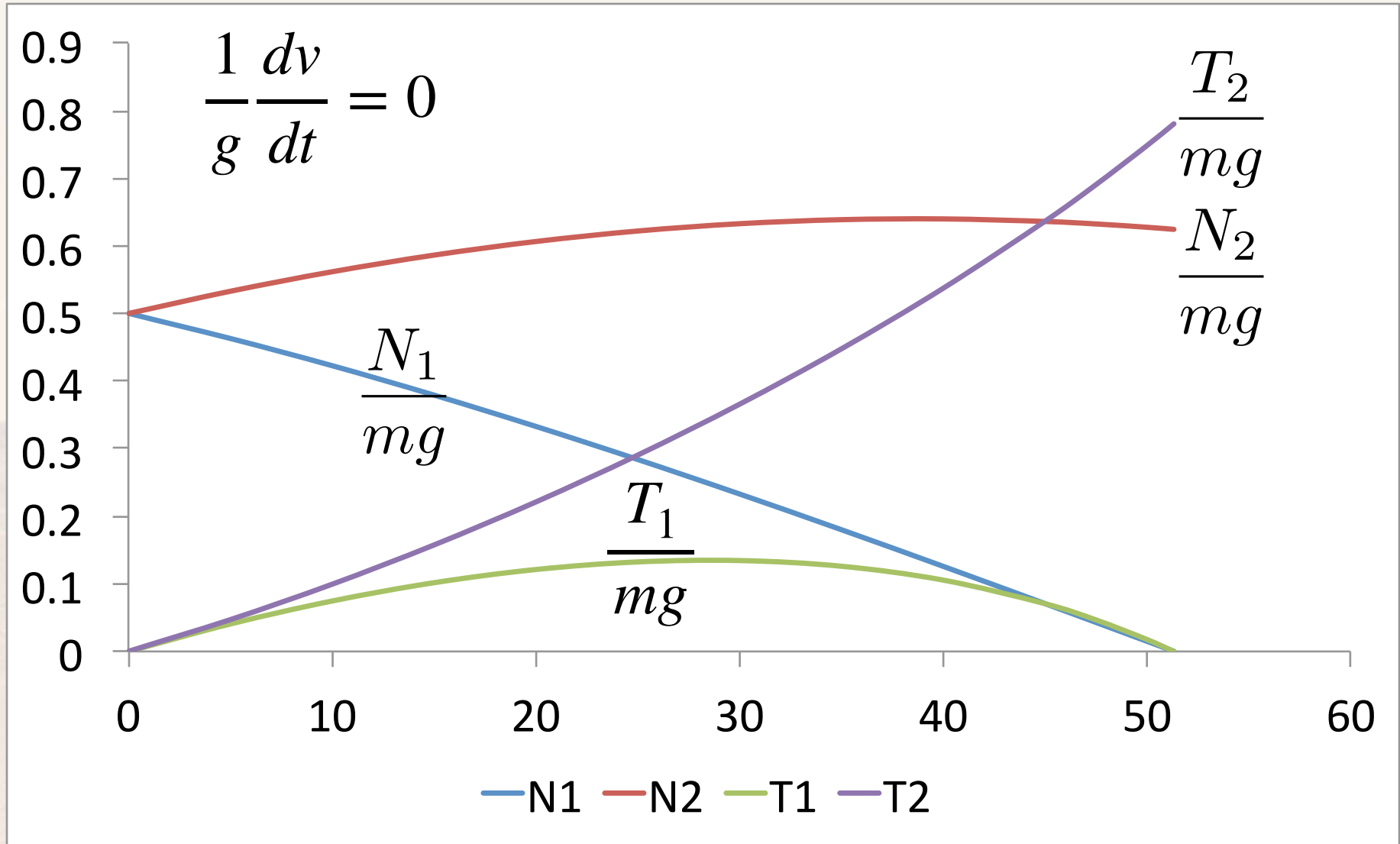
$$N_2 = mg \left[\frac{a}{\ell} \cos \theta + \left(\frac{h}{\ell} + \frac{r}{\ell} \right) \sin \theta + \frac{1}{g} \frac{dv}{dt} \right]$$

$$T_2 = \frac{N_2}{N_1 + N_2} \left(mg \sin \theta + m \frac{dv}{dt} \right)$$

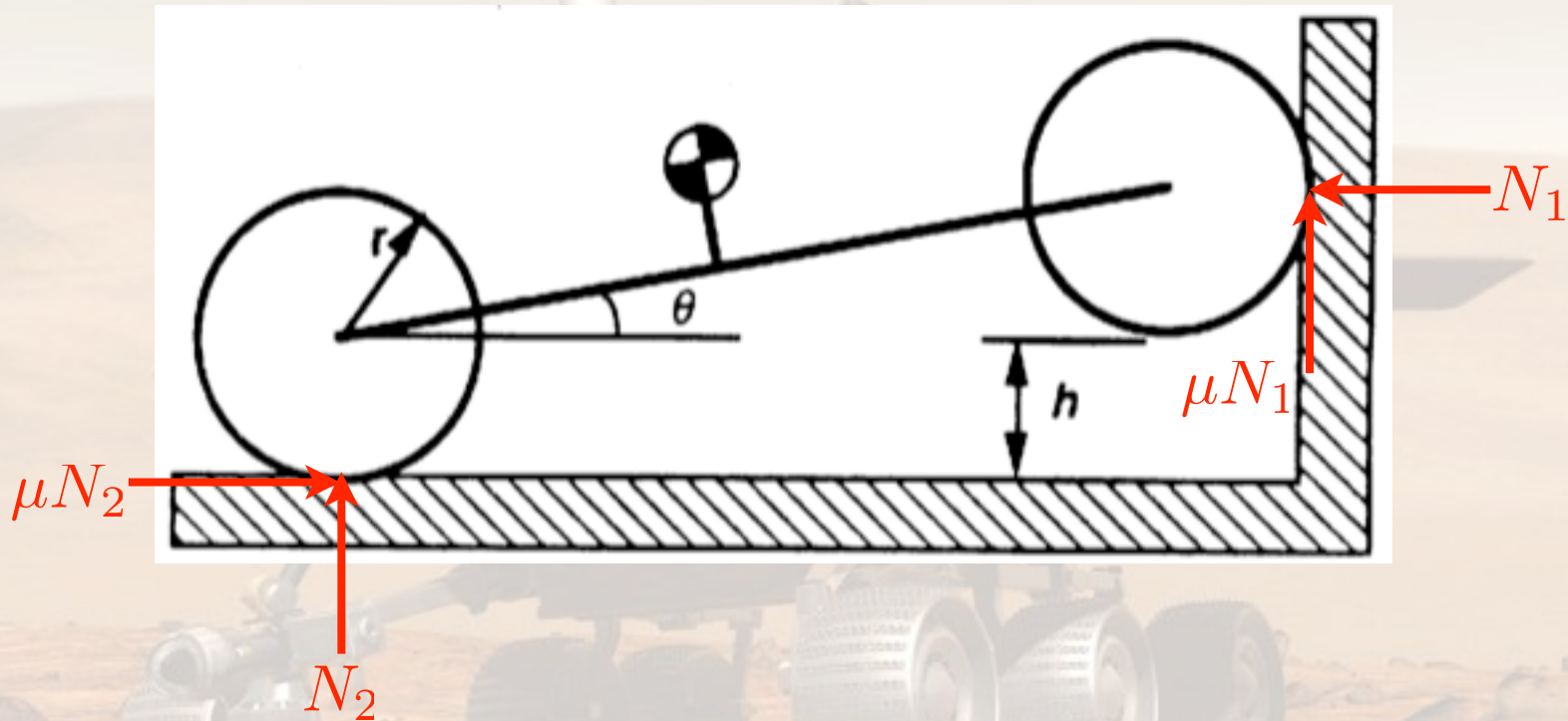
$$T_1 = \frac{N_1}{N_1 + N_2} \left(mg \sin \theta + m \frac{dv}{dt} \right)$$



Normal and Shear Wheel Force w/Slope



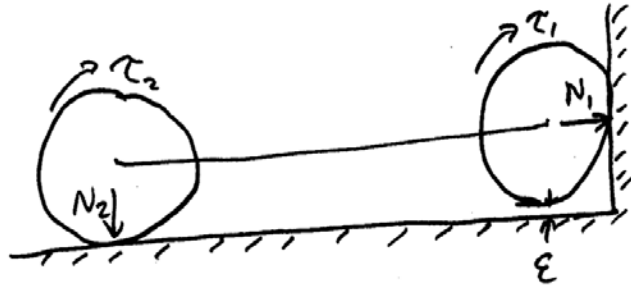
Four-Wheeled Vehicle Climbing a Wall



from Howard Eisen, "Scale and Computer Modeling of Wheeled Vehicles for Planetary Exploration" S.M. Thesis,
Department of Aeronautics and Astronautics, Massachusetts Institute of Technology, May, 1990



Wall Climbing



$$\sum F_{\text{vertical}} \Rightarrow N_2 + \mu N_1 = W$$

$$\sum F_{\text{horizontal}} \Rightarrow \mu N_2 = N_1$$

$$\sum M_{\text{rear axle}} \Rightarrow \mu N_2 r + \mu N_1 (l+r) = W(l-a)$$

$$N_2 + \mu^2 N_2 = W \Rightarrow N_2 = \frac{W}{1+\mu^2} \Rightarrow N_1 = \frac{\mu}{1+\mu^2} W$$

$$\frac{\mu}{1+\mu^2} W r + \frac{\mu^2}{1+\mu^2} W (l+r) = W(l-a)$$

$$(179) \mu^2 + r\mu - (l-a) = 0$$

$$\mu = \frac{-r \pm \sqrt{r^2 + 4(r+a)(l-a)}}{2(r+a)}$$

$$\text{Let } \alpha \equiv \frac{a}{r} \quad \lambda \equiv \frac{l}{r}$$

$$\mu = \frac{-1 \pm \sqrt{1 + 4(1+\alpha)(\lambda-\alpha)}}{2(1+\alpha)}$$

$$\text{Assume } \alpha = \frac{\lambda}{2}$$

$$\mu = -\frac{1}{2+\lambda} \pm \frac{\sqrt{1 + 4(1 + \frac{\lambda}{2})(\frac{\lambda}{2})}}{2+\lambda} \quad \left. \vphantom{\mu} \right\} \rightarrow$$

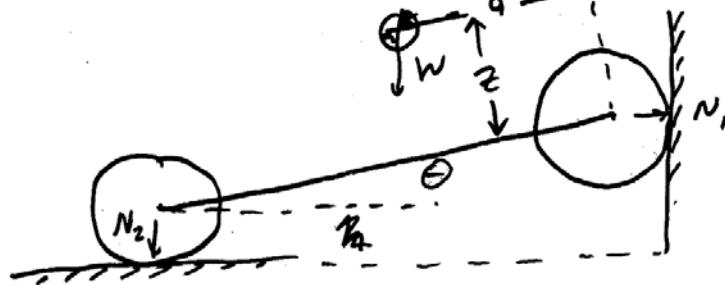
$$\begin{aligned} & 1 + 4\frac{\lambda}{2} + 4\frac{\lambda^2}{4} \\ &= 1 + 2\lambda + \lambda^2 \\ &= (\lambda + 1)^2 \end{aligned}$$

$$= -\frac{1 \pm (\lambda + 1)}{2 + \lambda} = \frac{\lambda}{2 + \lambda}, -1$$

$$\lambda \rightarrow 0 \quad \mu_{\text{limit}} \rightarrow 0 \quad \lambda \rightarrow \infty \quad \mu_{\text{limit}} \rightarrow 1$$

Shorter is better!

A Short Time Later...



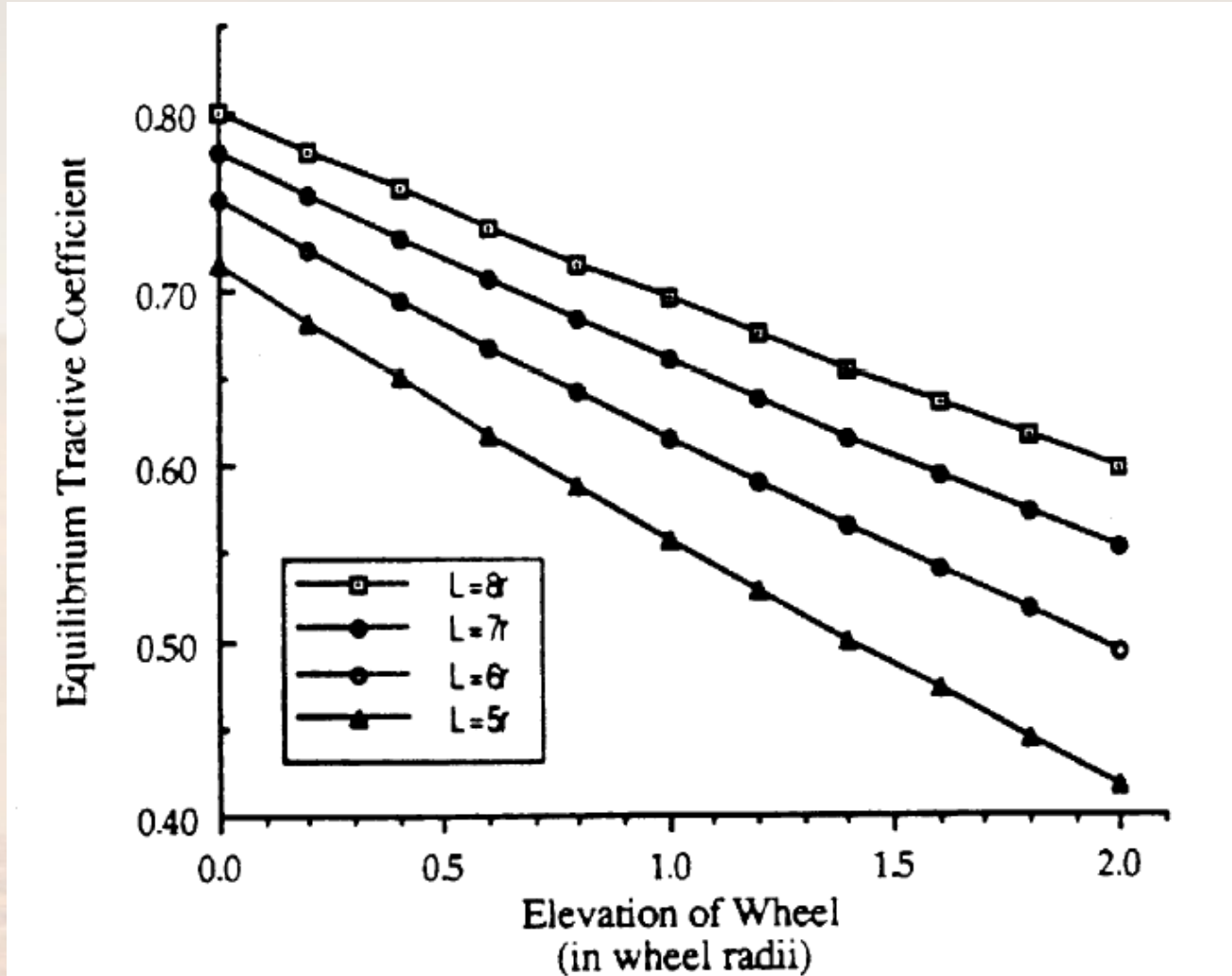
$$h = l \sin \theta$$

$$\left. \begin{aligned} \Sigma F_{\text{vert}} &\Rightarrow \mu N_1 + N_2 = W \\ \Sigma F_{\text{horiz}} &\Rightarrow \mu N_2 = N_1 \end{aligned} \right\} \begin{aligned} N_1 &= \frac{\mu}{1+\mu^2} W \\ N_2 &= \frac{W}{1+\mu^2} \end{aligned}$$

$$\Sigma M_{\text{rear}} \Rightarrow \mu N_2 r + N_1 l \sin \theta + \mu N_1 (r + l \cos \theta) = W [(l-a) \cos \theta - z \sin \theta]$$

As θ increases, effect of W decreases and effect of N_1 increases \Rightarrow hardest point of the climb is at the start!

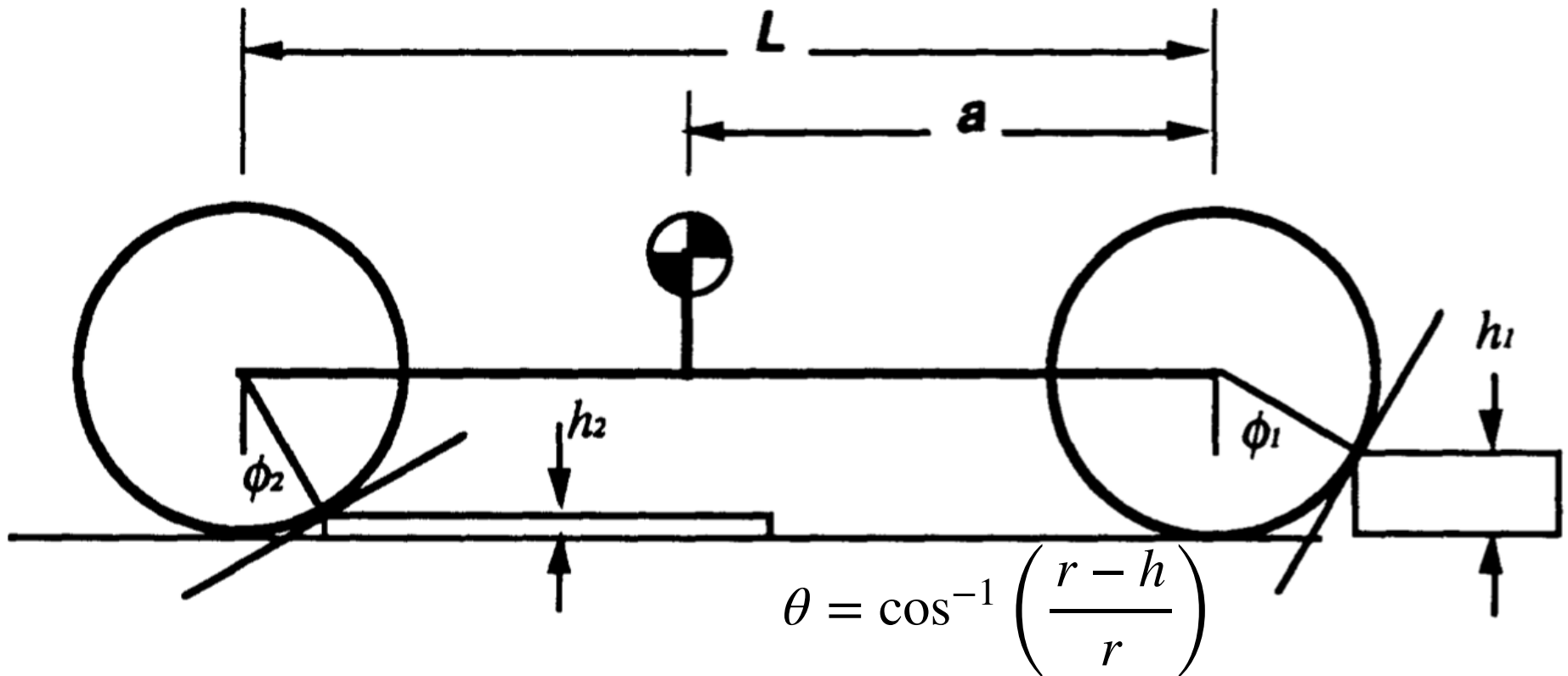
Required Traction for Wall Climbing



from Howard Eisen, "Scale and Computer Modeling of Wheeled Vehicles for Planetary Exploration" S.M. Thesis, Department of Aeronautics and Astronautics, Massachusetts Institute of Technology, May, 1990



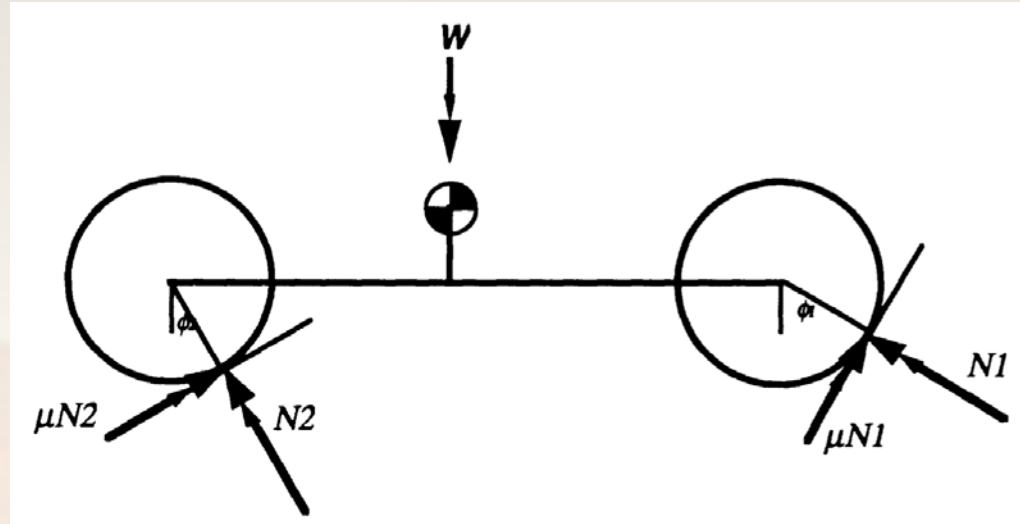
Wheel Interaction with Slope



from Howard Eisen, "Scale and Computer Modeling of Wheeled Vehicles for Planetary Exploration" S.M. Thesis, Department of Aeronautics and Astronautics, Massachusetts Institute of Technology, May, 1990



Equations for Slopes under Wheels



Sum of Horizontal forces:

$$\mu N_2 \sin \phi_2 + N_2 \cos \phi_2 + \mu N_1 \sin \phi_1 + N_1 \cos \phi_1 - W = 0$$

Sum of vertical forces:

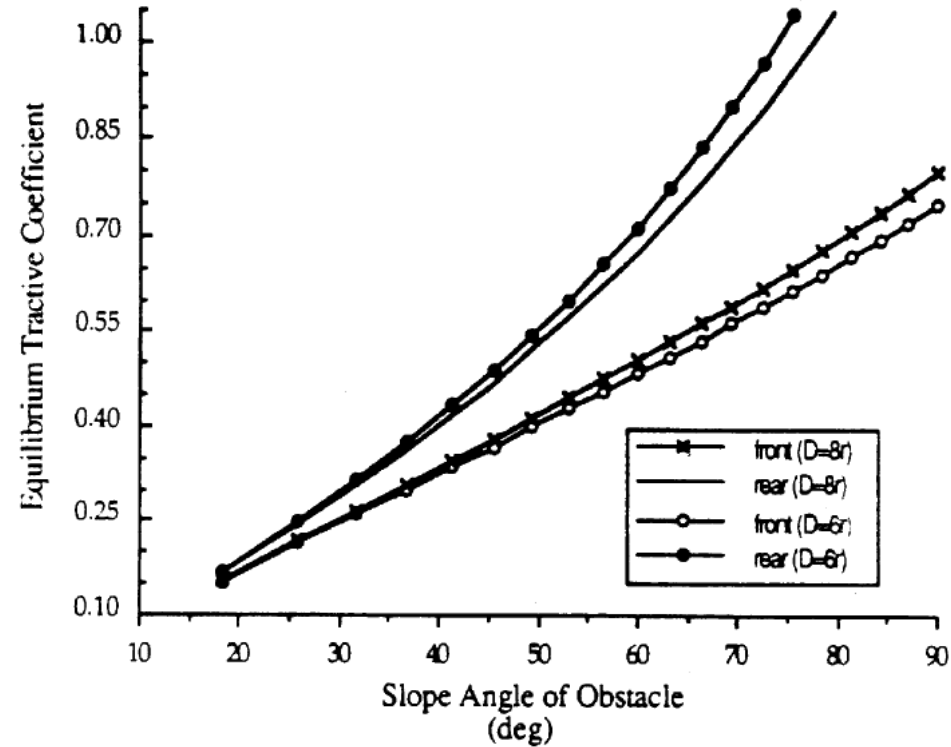
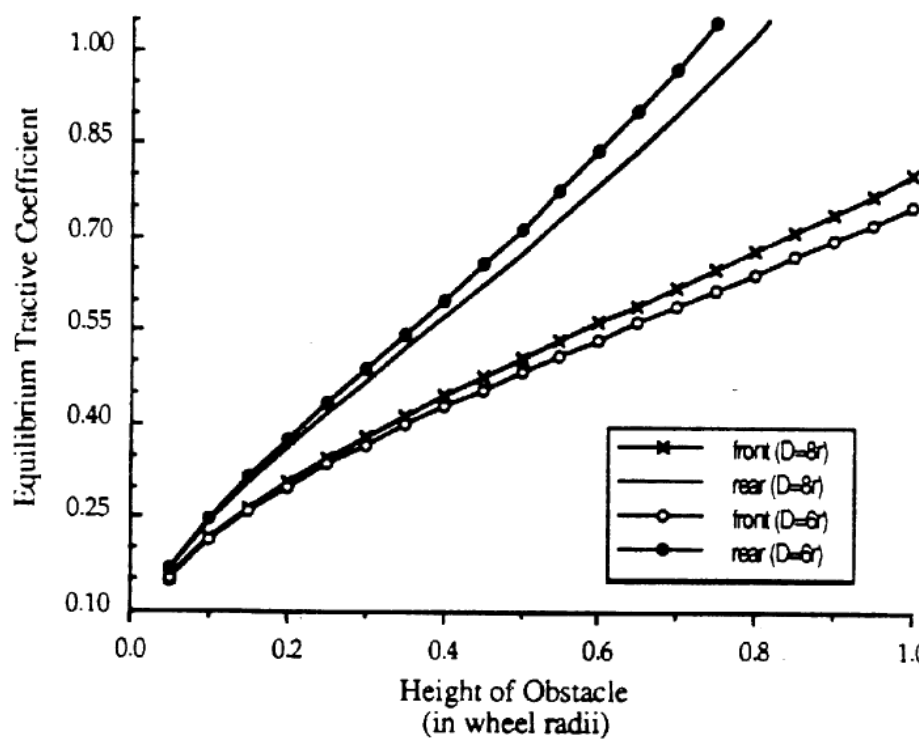
$$\mu N_2 \cos \phi_2 - N_2 \sin \phi_2 + \mu N_1 \cos \phi_1 - N_1 \sin \phi_1 = 0$$

Sum of forces around the rear axle:

$$\left(\mu N_2 r - W(L - a) + N_1 L \cos \phi_1 + \mu N_1 (r + L \sin \phi_1) \right) = 0$$



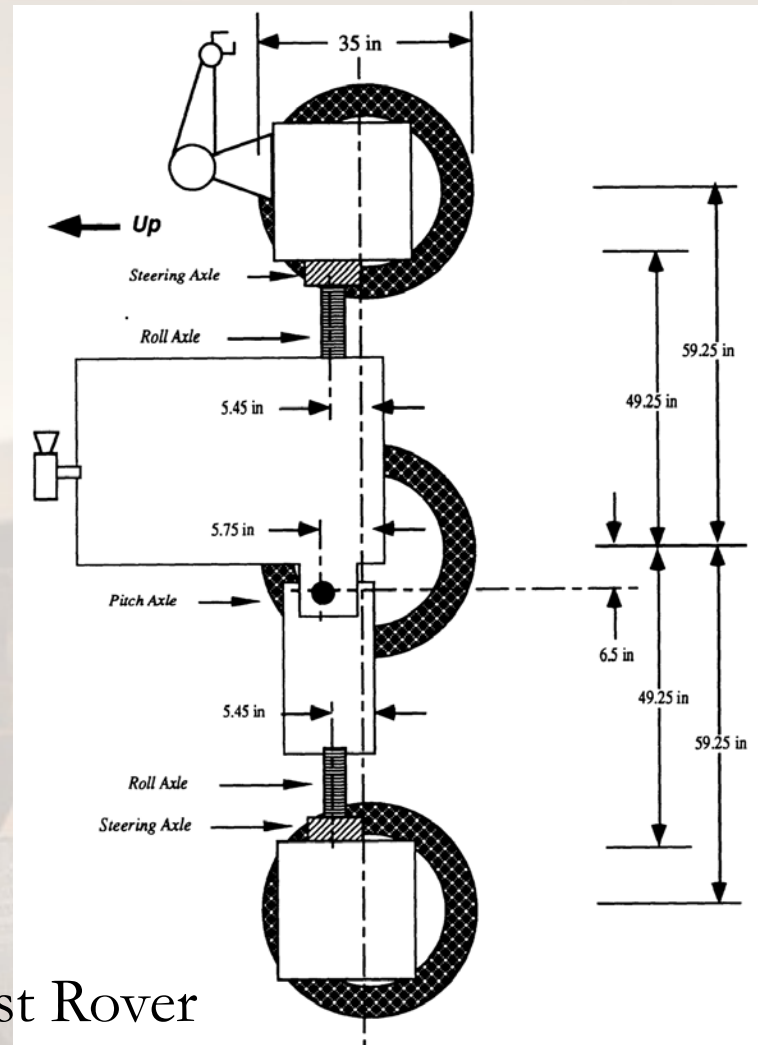
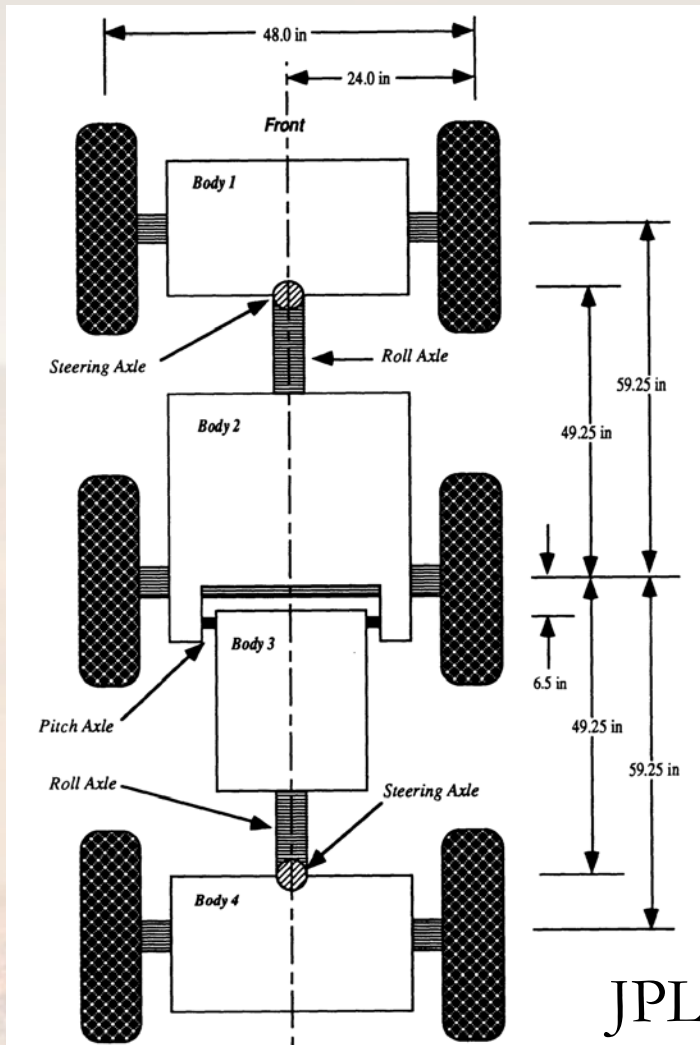
Bump/Slope Traction Requirements



from Howard Eisen, "Scale and Computer Modeling of Wheeled Vehicles for Planetary Exploration" S.M. Thesis, Department of Aeronautics and Astronautics, Massachusetts Institute of Technology, May, 1990



Six-Wheel Articulated Body Rover

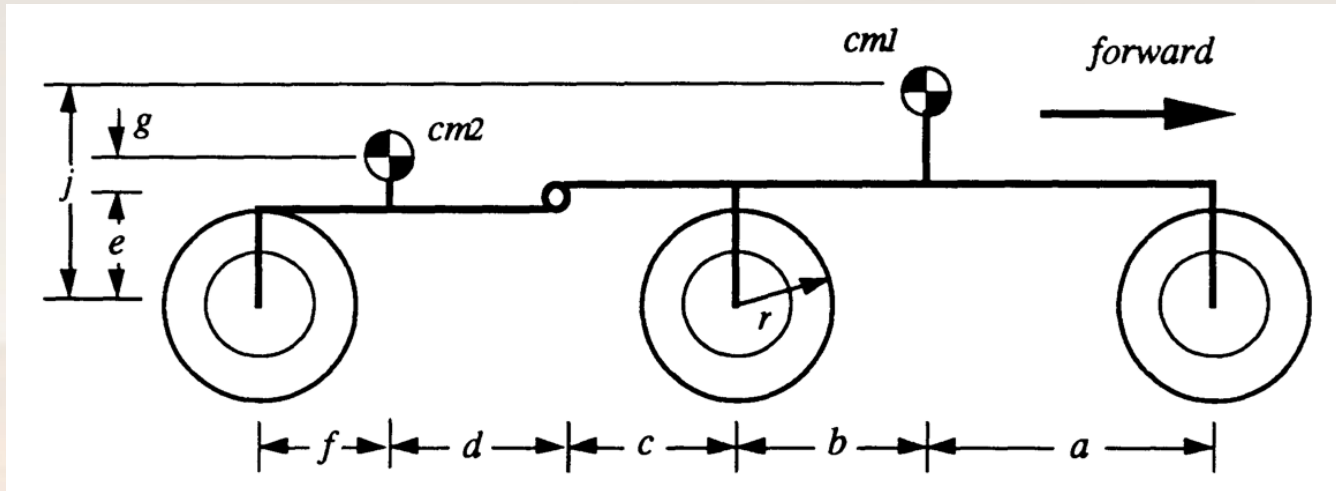


JPL Navtest Rover

from Howard Eisen, "Scale and Computer Modeling of Wheeled Vehicles for Planetary Exploration" S.M. Thesis, Department of Aeronautics and Astronautics, Massachusetts Institute of Technology, May, 1990



Model of Six-Wheel Vehicle



Sum of vertical forces:

$$N_3 + \mu N_2 + N_1 - Wf - Wb = 0$$

Sum of horizontal forces:

$$\mu N_3 - N_2 + \mu N_1 = 0$$

Sum of moments for front body around pitch axis

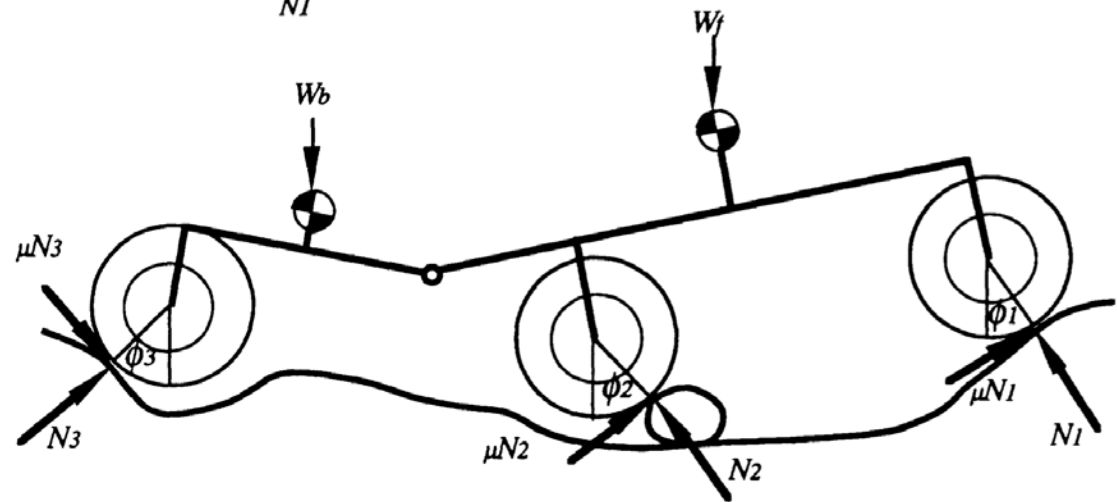
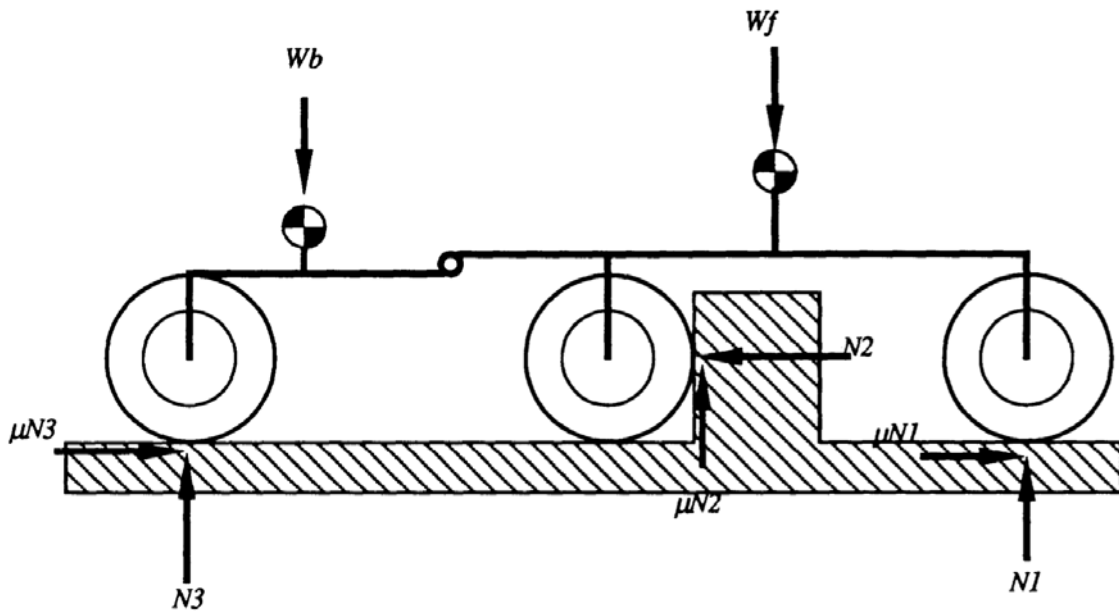
$$\mu N_1(r + e) + N_1(a + b + c) + \mu N_2(r + c) + -N_2e - Wf(b + c) = 0$$

Sum of moments for rear body around pitch axis

$$Wbd + \mu N_3(r + e) - N_3(d + f) = 0$$



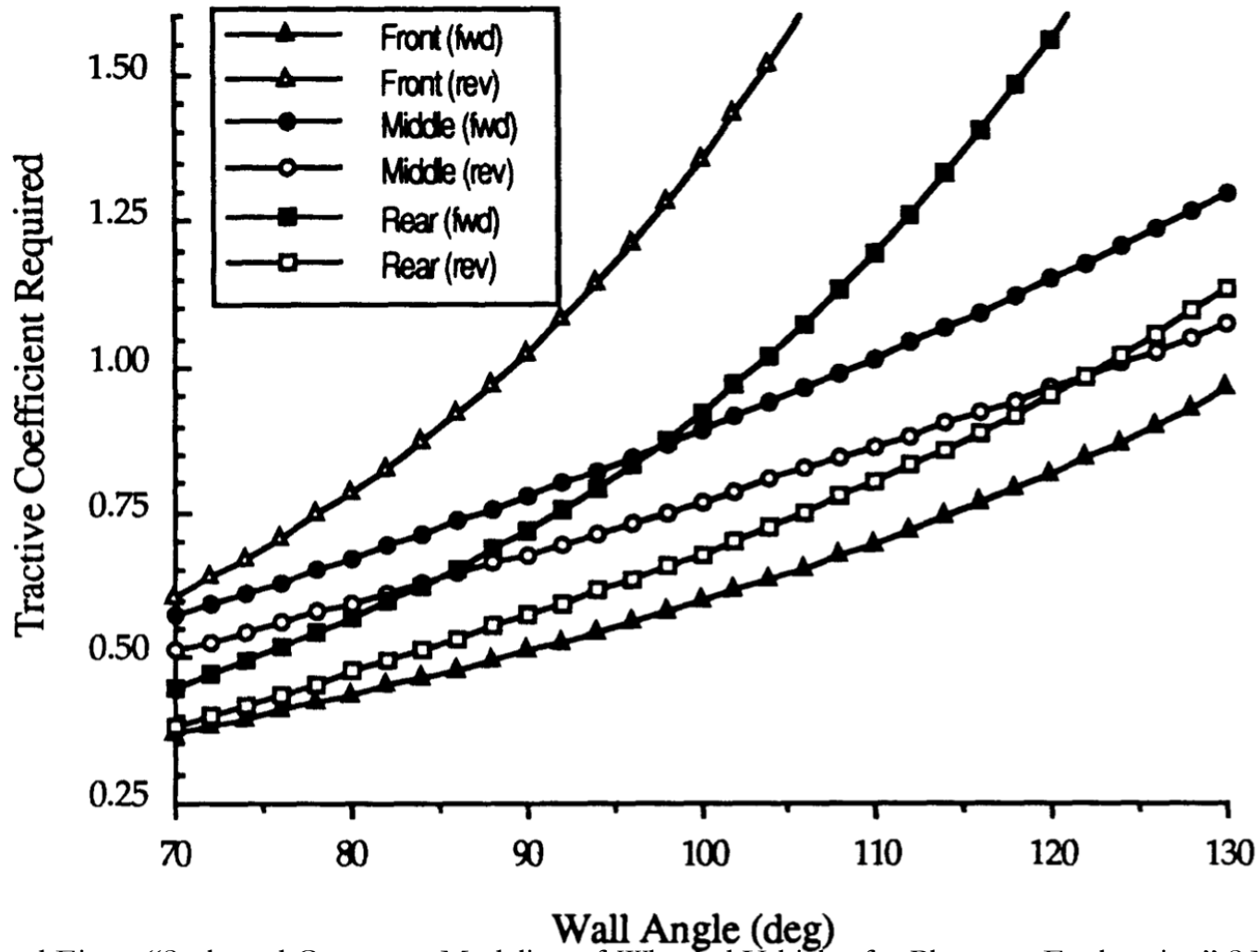
Navtest Rover with Walls and Slopes



from Howard Eisen, "Scale and Computer Modeling of Wheeled Vehicles for Planetary Exploration" S.M. Thesis, Department of Aeronautics and Astronautics, Massachusetts Institute of Technology, May, 1990



Six-Wheel Rover, Slope Climbing

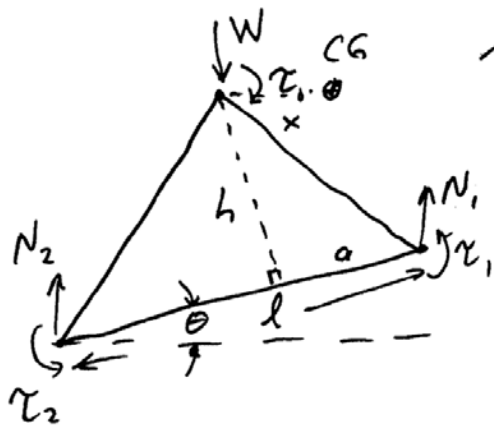


from Howard Eisen, "Scale and Computer Modeling of Wheeled Vehicles for Planetary Exploration" S.M. Thesis, Department of Aeronautics and Astronautics, Massachusetts Institute of Technology, May, 1990



Four-Wheel Rocker Suspension





$$\tau_0 = xW$$

Planar Rocker Analysis

$$\Sigma \text{ Forces: } N_1 + N_2 = W$$

$$\Sigma \text{ Moment (rear axle)}$$

$$\tau_1 + \tau_2 + N_1 l \cos \theta = \tau_0 + W[(l-a) \cos \theta - h \sin \theta]$$

$$N_1 l \cos \theta = \tau_0 - \tau_1 - \tau_2 + W[(l-a) \cos \theta - h \sin \theta]$$

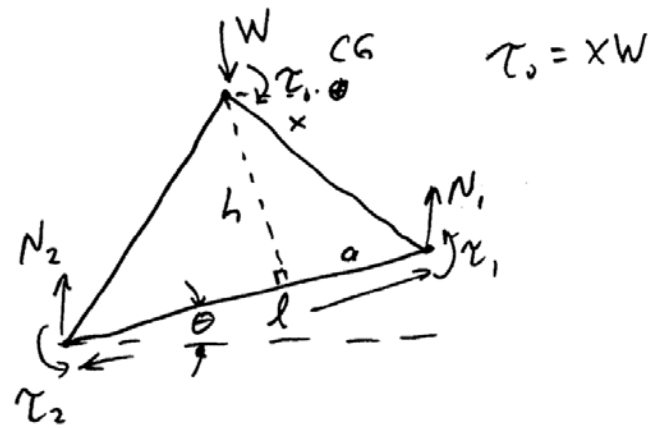
$$N_1 = \frac{\tau_0 - \tau_1 - \tau_2}{l \cos \theta} + W \left(\frac{l-a}{l} - \frac{h}{l} \tan \theta \right)$$


$$N_2 = W - N_1 = W \left(1 - \frac{l-a}{l} + \frac{h}{l} \tan \theta \right) - \frac{\tau_0 - \tau_1 - \tau_2}{l \cos \theta}$$

$$N_2 = W \left(\frac{a}{l} + \frac{h}{l} \tan \theta \right) + \frac{\tau_1 + \tau_2 - \tau_0}{l \cos \theta}$$

Non-dimensionalize

$$\tau_o = Wx \Rightarrow \frac{\tau_o}{Wl} = \frac{x}{l}$$





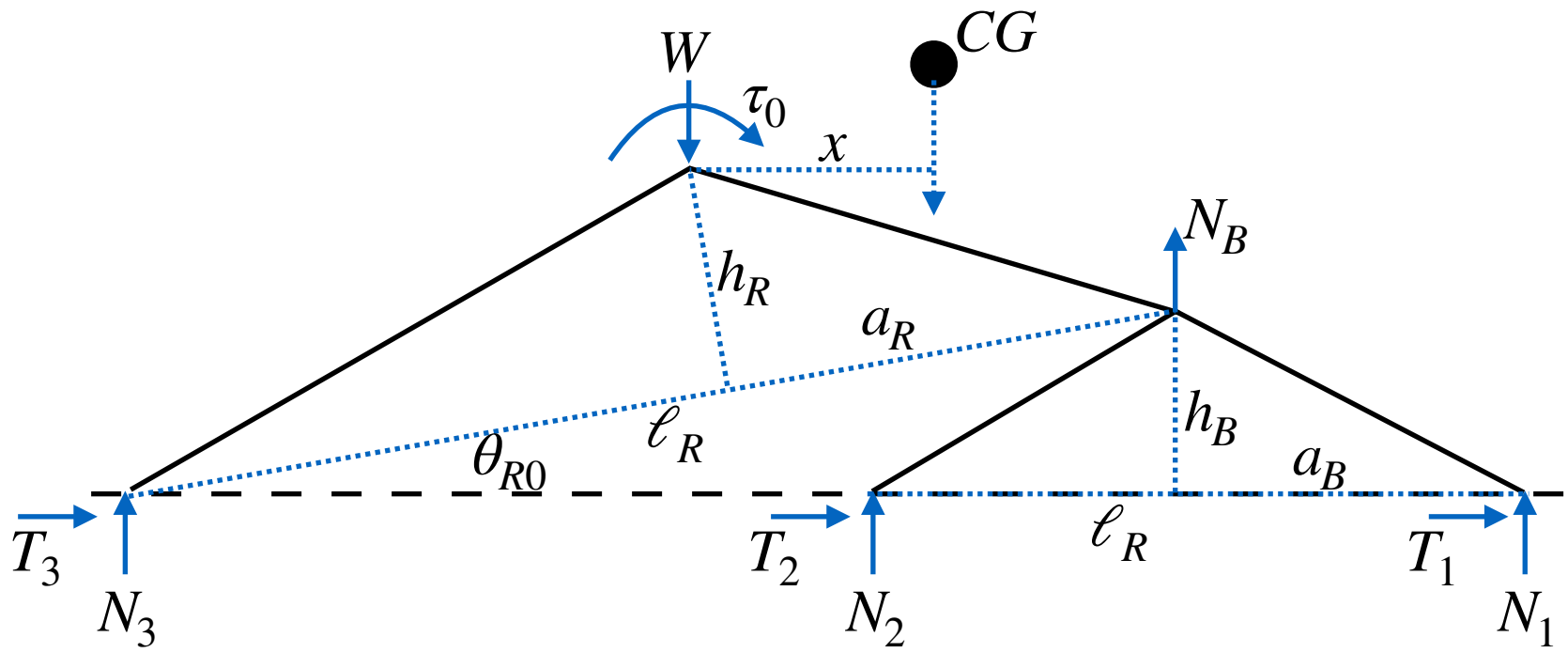
$$\tau_{\text{wheel}} = Tr \Rightarrow \frac{\tau_{\text{wheel}}}{Wl} = \frac{T}{W} \frac{r}{l}$$

$$\frac{N_1}{W} = 1 - \frac{a}{l} - \frac{h}{l} \tan \theta + \left(\frac{x}{l} - \frac{T_1}{W} \frac{r_1}{l} - \frac{T_2}{W} \frac{r_2}{l} \right) \frac{1}{\cos \theta}$$

$$\frac{N_2}{W} = \frac{a}{l} + \frac{h}{l} \tan \theta + \left(\frac{T_1}{W} \frac{r_1}{l} + \frac{T_2}{W} \frac{r_2}{l} - \frac{x}{l} \right) \frac{1}{\cos \theta}$$

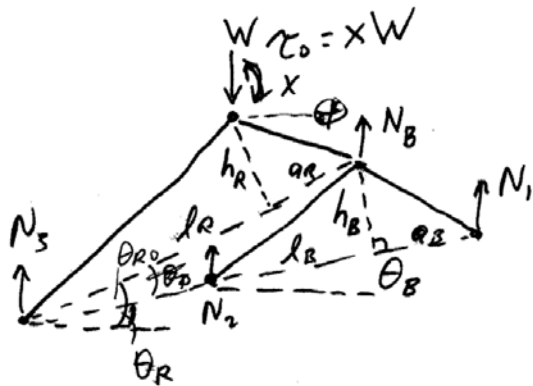
Six-Wheel Rocker-Bogey Suspension





Kinematics of Planar Rocker-Bogey

Planar Rocker - Bogey Analysis



For bogey,

$$N_1 = N_B \left(\frac{l_B \cdot a_B}{l_B} - \frac{h_B}{l_B} \tan \theta_B \right) - \frac{\tau_1 + \tau_2}{l_B \cos \theta_B}$$

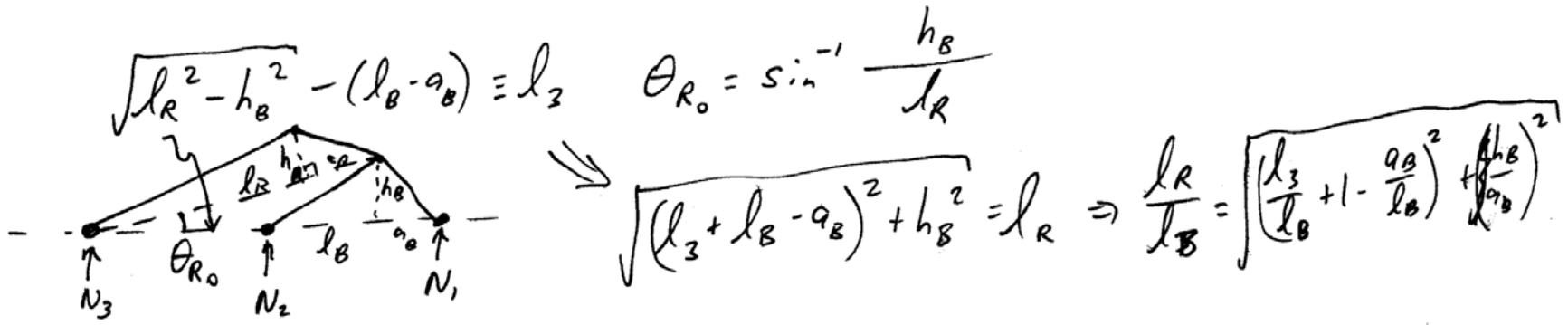
$$N_2 = N_B \left(\frac{a_B}{l_B} + \frac{h_B}{l_B} \tan \theta_B \right) + \frac{\tau_1 + \tau_2}{l_B \cos \theta_B}$$

($\tau_B = 0$ always)

For rocker,

$$N_B = \frac{\tau_0 - \tau_3}{l_R \cos(\theta_R + \theta_{R_0})} + W \left(\frac{l_R - a_R}{l_R} - \frac{h_R}{l_R} \tan(\theta_R + \theta_{R_0}) \right)$$

$$N_3 = W \left(\frac{a_R}{l_R} + \frac{h_R}{l_R} \tan(\theta_R + \theta_{R_0}) \right) + \frac{\tau_3 - \tau_0}{l_R \cos(\theta_R + \theta_{R_0})}$$



Normalize by W and l_B

$$\frac{N_3}{W} = \frac{a_R}{l_R} + \frac{h_R}{l_R} \tan(\theta_R + \theta_{R0}) + \left(\frac{T_3}{W} \frac{r_3}{l_B} \frac{l_B}{l_R} - \frac{x}{l_B} \frac{l_B}{l_R} \right) \frac{1}{\cos(\theta_R + \theta_{R0})}$$

$$\frac{N_B}{W} = 1 - \frac{a_R}{l_R} - \frac{h_R}{l_R} \tan(\theta_R + \theta_{R0}) + \left(\frac{x}{l_B} \frac{l_B}{l_R} - \frac{T_3}{W} \frac{r_3}{l_B} \frac{l_B}{l_R} \right) \frac{1}{\cos(\theta_R + \theta_{R0})}$$

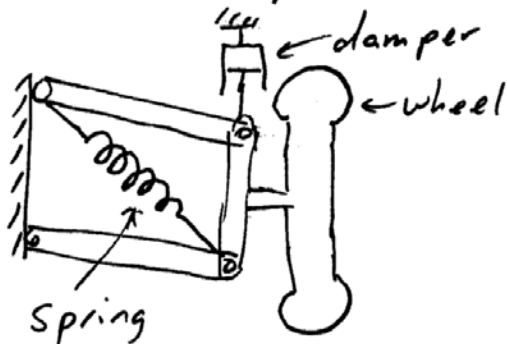
$$\frac{N_1}{W} = \frac{N_B}{W} \left(1 - \frac{a_B}{l_B} - \frac{h_B}{l_B} \tan \theta_B \right) - \left(\frac{T_1}{W} \frac{r_1}{l_B} + \frac{T_2}{W} \frac{r_2}{l_B} \right) \frac{1}{\cos \theta_B}$$

$$\frac{N_2}{W} = \frac{N_B}{W} \left(\frac{a_B}{l_B} + \frac{h_B}{l_B} \tan \theta_B \right) + \left(\frac{T_1}{W} \frac{r_1}{l_B} + \frac{T_2}{W} \frac{r_2}{l_B} \right) \frac{1}{\cos \theta_B}$$

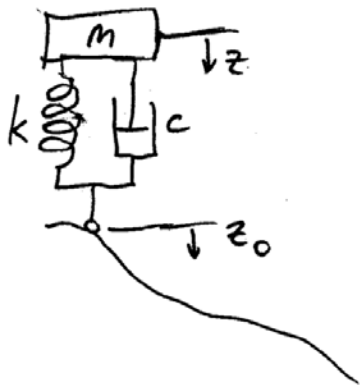
Suspension Systems

- Current planetary rovers (e.g., MER, MSL) have little or no shock absorption

- Notional car suspension



Analyze one wheel first



$$m\ddot{z} + c\dot{z} + kz = c\dot{z}_0 + kz_0$$

Undamped force-free solution

$$m\ddot{z} + kz = 0 \quad z = Z_1 \cos \omega_n t$$

$$-m\omega_n^2 + k = 0 \quad \ddot{z} = -Z_1 \omega_n^2 \cos \omega_n t$$

$$\omega_n = \sqrt{\frac{k}{m}}$$

Rover Example

$$M_{TOT} = 500 \text{ kg} \Rightarrow \text{wheel } m = \frac{M_{TOT}}{4} = 125 \text{ kg}$$

d = deflection of suspension (at rest) $\sim 0.1 \text{ m}$

$$k = \frac{F}{d} = \frac{mg}{d}$$

$$\omega_n = \sqrt{\frac{k}{m}}$$

$$f_n = \frac{\omega_n}{2\pi}$$

	<u>Earth</u>	<u>Moon</u>
k	12,250 N/m	2000 N/m
ω_n	9.9 rad/sec	4 rad/sec
f_n	1.6 Hz	0.64 Hz

l_{crit} = critical distance
between bumps

$$= \frac{V}{f_n}$$

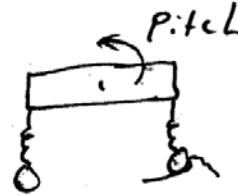
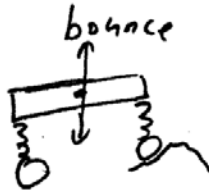
@10kph
(2.8 m/sec)

l_{crit} 1.8m

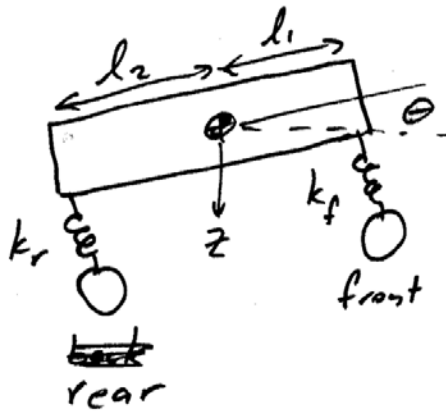
4.3m

Multiwheel Analysis

Responses to two-wheels hitting a bump



We can excite both of these modes



Equation of Motion
(assuming no damping)

Bounce: $m \ddot{z} + k_f(z - l_1 \theta) + k_r(z + l_2 \theta) = 0$
(small angles)

Pitch: $I_y \ddot{\theta} + k_f l_1 (z - l_1 \theta) + k_r l_2 (z + l_2 \theta) = 0$

let $I_y = m r_y^2$ $r_y = \text{radius of gyration}$

Solve this set of coupled differential eqns

let $D_1 \equiv \frac{k_f + k_r}{m}$ $D_2 \equiv \frac{k_r l_2 - k_f l_1}{m}$ $D_3 \equiv \frac{k_f l_1^2 + k_r l_2^2}{I_y}$

Rewrite in terms of D_1, D_2, D_3

$$\ddot{z} + D_1 z + D_2 \theta = 0$$

$D_2 =$ coupling coefficient

$$\ddot{\theta} + D_3 \theta + \frac{D_2}{r_y^2} z = 0$$

Equations are independent

if $D_2 = 0 \Rightarrow k_f l_1 = k_r l_2$

If $D_2 = 0$, force @ CG only produces bounce $\omega_{nz} = \sqrt{D_1}$

force elsewhere produces pitch $\omega_{n\theta} = \sqrt{D_3}$

Assume $D_2 \neq 0$

$$z = Z \cos \omega_n t$$

$$\theta = \Theta \cos \omega_n t$$

$$\left. \begin{aligned} (D_1 - \omega_n^2) Z + D_2 \Theta &= 0 \\ \frac{D_2}{r_y^2} Z + (D_3 - \omega_n^2) \Theta &= 0 \end{aligned} \right\} \begin{vmatrix} D_1 - \omega_n^2 & D_2 \\ \frac{D_2}{r_y^2} & D_3 - \omega_n^2 \end{vmatrix} = 0$$

$$\omega_n^4 - (D_1 + D_3) \omega_n^2 + \left(D_1 D_3 - \frac{D_2^2}{r_y^2} \right) \omega_n^2 - (D_1 + D_3) \omega_n^2 + \left(D_1 D_3 - \frac{D_2^2}{r_y^2} \right) = 0$$

$$\omega_{n1}^2 = \frac{D_1 + D_3}{2} \pm \frac{1}{2} \sqrt{(D_1 + D_3)^2 - 4 \left(D_1 D_3 - \frac{D_2^2}{r_y^2} \right)}$$

∴ (Insert algebra here)

$$\omega_{n1}^2 = \frac{D_1 + D_3}{2} - \sqrt{\frac{1}{4} (D_1 - D_3)^2 + \frac{D_2^2}{r_y^2}}$$

$$\omega_{n2}^2 = \frac{D_1 + D_3}{2} + \sqrt{\frac{1}{4} (D_1 - D_3)^2 + \frac{D_2^2}{r_y^2}}$$

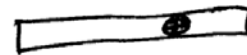
Example: $k_f = k_r = 2000 \text{ N/m}$ (moon)

$$l_1 = 1 \text{ m} \quad l_2 = 2 \text{ m}$$

$$D_1 = \frac{4000}{500} = 8 \frac{\text{N/m}}{\text{kg}} = \left\langle \frac{1}{\text{sec}^2} \right\rangle$$

$$D_2 = \frac{2000 \frac{\text{N}}{\text{m}} (2 \text{ m}) - 2000 \frac{\text{N}}{\text{m}} (1 \text{ m})}{500 \text{ kg}} = 4 \frac{\text{N}}{\text{kg}} = 4 \frac{\text{m}}{\text{sec}^2}$$

$$D_3 = \frac{2000 \frac{\text{N}}{\text{m}} (1 \text{ m})^2 + 2000 \frac{\text{N}}{\text{m}} (2 \text{ m})^2}{375 \text{ kg m}^2} = 26.7 \frac{\text{N m}}{\text{kg m}^2} = \left\langle \frac{1}{\text{sec}^2} \right\rangle$$



$$I = \frac{ml^2}{12} \Rightarrow I_y = 375 \text{ kg m}^2$$

$$r_y = 0.75 \text{ m}$$

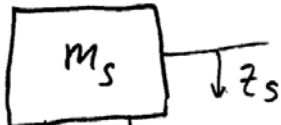
$$\omega_n^2 = 17.33 \pm 10.43$$

$$\omega_{n_1} = 2.63 \text{ rad/sec} \Rightarrow 0.42 \text{ Hz}$$

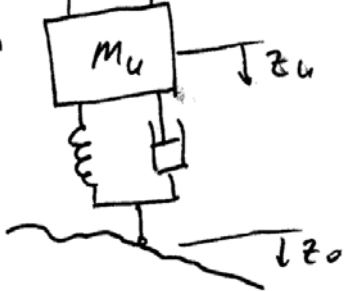
$$\omega_{n_2} = 5.67 \text{ rad/sec} \Rightarrow 0.84 \text{ Hz}$$

Add in Tire Mass & Stiffness

"Sprung"
mass



"unsprung"
mass



Sprung mass

$$m_s \ddot{z}_s + C_s (\dot{z}_s - \dot{z}_u) + k_s (z_s - z_u) = 0$$

unsprung mass

$$m_u \ddot{z}_u + C_g (\dot{z}_u - \dot{z}_s) + k_s (z_u - z_s) + C_u \dot{z}_u + k_u z_u = F(t) = C_u \dot{z}_0 + k_u z_0$$

Undamped Force-Free Solutions

$$m_s \ddot{z}_s + k_s (z_s - z_u) = 0$$

$$m_u \ddot{z}_u + k_s (z_u - z_s) + k_u z_u = 0$$

$$z_s = Z_s \cos \omega_n t \quad z_u = Z_u \cos \omega_n t$$

$$\begin{vmatrix} k_s - m_s \omega_n^2 & -k_u \\ -k_s & k_s + k_u - m_u \omega_n^2 \end{vmatrix} = 0$$

$$\omega_n^4 (m_u m_s) + \omega_n^2 (-m_s k_s - m_s k_u - m_u k_s) + k_s k_u = 0$$

$$\omega_{n1} = \frac{+B \pm \sqrt{B^2 - 4AC}}{2A}$$

$$\omega_{n2} = \frac{B + \sqrt{B^2 - 4AC}}{2A}$$

$$A = m_u m_s \quad B = m_s (k_s + k_u) + m_u k_s \quad C = k_s k_u$$

Example:

$$m_s = 100 \text{ kg} \quad m_u = 25 \text{ kg}$$

$$k_s = 2000 \text{ N/m} \quad k_u = 10,000 \text{ N/m}$$

$$A = 2500 \text{ kg}^2$$

$$B = 1.25 \times 10^6 \text{ kg}^2/\text{sec}^2$$

$$C = 2 \times 10^7 \text{ N}^2/\text{m}^2$$

$$\omega_{n1} = 4.76 \frac{\text{rad}}{\text{sec}} \Rightarrow 0.8 \text{ Hz} \leftarrow \text{suspension frequency}$$

$$\omega_{n2} = 28.8 \frac{\text{rad}}{\text{sec}} \Rightarrow 3.5 \text{ Hz} \leftarrow \text{wheel } \overset{\text{stiffness}}{\text{response}} \text{ frequency}$$

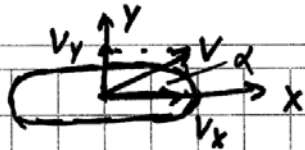
Wheel-Soil Interaction in Turn



Slip ratios

$$s = \frac{r\omega - v_x}{r\omega} \quad \text{for driving: } |r\omega| > |v_x|$$

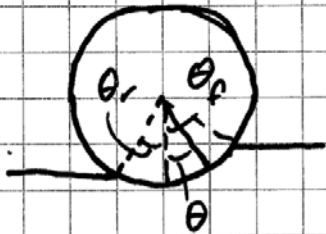
$$s = \frac{r\omega - v_x}{v_x} \quad \text{for braking: } |r\omega| < |v_x|$$



$$-1 \leq s \leq 1$$

$$\text{Slip angle } \alpha = \tan^{-1}\left(\frac{v_y}{v_x}\right)$$

Wheel sinkage $P(z) = \left(\frac{k_c}{b} + k_\phi\right) z^n$



$$z(\theta) = r(\cos \theta - \cos \theta_s)$$

for static sinkage, $\theta_f = \theta_r = \theta_s$

$$z(\theta) = r(\cos \theta - \cos \theta_s)$$

$$P(\theta) = \left(\frac{k_c}{b} + k_\phi\right) r^n (\cos \theta - \cos \theta_s)^n$$

Given weight W on wheel,

$$W = \int_{-\theta_s}^{\theta_s} P(\theta) b r \cos \theta d\theta = r^{n+1} (k_c + k_\phi b) \int_{-\theta_s}^{\theta_s} (\cos \theta - \cos \theta_s)^n \cos \theta d\theta$$

Static sinkage $z_s = r(1 - \cos \theta_s)$

Soil entry angle $\theta_f = \cos^{-1} \left(1 - \frac{z_{max}}{r} \right)$

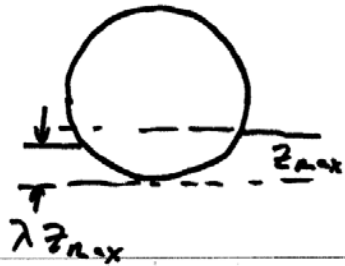
exit angle $\theta_r = \cos^{-1} \left(1 - \lambda \frac{z_{max}}{r} \right)$

λ = wheel sinkage ratio

= 0 when no soil restitution occurs

= 1 if soil completely rebounds (static sinkage)

can be > 1 if soil is transported to back of wheel



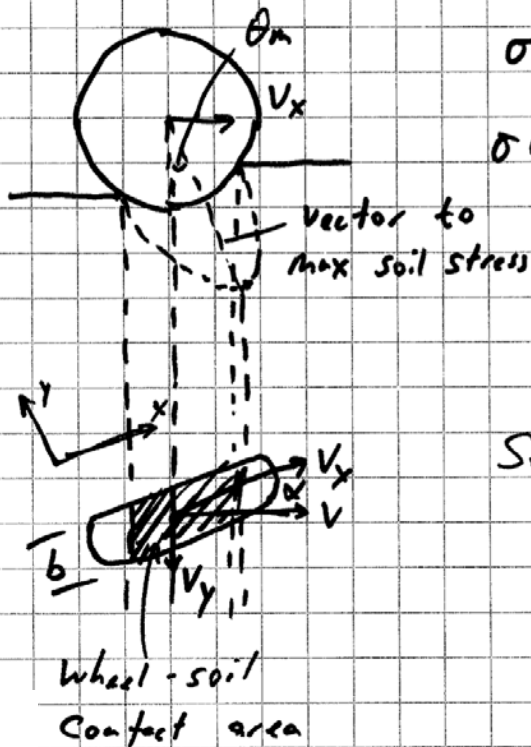
Wheel Stress Distribution

$$\sigma(\theta) = r^n \left(\frac{k_c}{b} + k_\phi \right) \left[\cos \theta - \cos \theta_f \right]^n \quad \theta_m \leq \theta < \theta_f$$

$$\sigma(\theta) = r^n \left(\frac{k_c}{b} + k_\phi \right) \left[\cos \left\{ \theta_f - \frac{\theta - \theta_r}{\theta_m - \theta_r} (\theta_f - \theta_m) \right\} - \cos \theta_f \right]^n \quad \theta_r < \theta \leq \theta_m$$

$$\theta_m = (q_0 + q_1 s) \theta_f$$

$$q_0 \sim 0.4 \quad 0 \leq q_1 \leq 0.3$$



Soil Shear Stresses

$$\tau_x(\theta) = (c + \sigma(\theta) \tan \phi) \left[1 - e^{-j_x(\theta)/k_x} \right]$$

$$\tau_y(\theta) = (c + \sigma(\theta) \tan \phi) \left[1 - e^{-j_y(\theta)/k_y} \right]$$

Soil Deformations

$$j_x(\theta) = r[\theta_f - \theta - (1-s)(\sin\theta_f - \sin\theta)]$$

$$j_y(\theta) = r(1-s)(\theta_f - \theta) \tan\alpha$$

Resolve into orthogonal forces

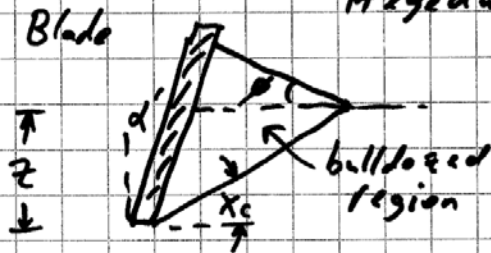
Drawbar pull $F_x = rb \int_{\theta_r}^{\theta_f} \{ \tau_x(\theta) \cos\theta - \sigma(\theta) \sin\theta \} d\theta$

Side force $F_y = \int_{\theta_r}^{\theta_f} \{ \underbrace{F_u}_{F_u} \tau_y(\theta) + \underbrace{F_s}_{F_s} R_b (r - z(\theta) \cos\theta) \} d\theta$

F_u are side loads generated under the wheel

F_s are side loads generated by bulldozing by side of wheel

Hagedorn bulldozing theory



$$R_b(z) = D_1 \left[c z(\theta) + D_2 \frac{\rho z^2(\theta)}{2} \right]$$

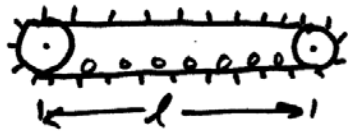
$$D_1(X_c, \phi) = \cot X_c + \tan(X_c + \phi)$$

$$D_2(X_c, \phi) = \cot X_c + \cot^2 X_c / \cot \phi$$

Bekker destruction angle $X_c = \frac{\pi}{4} - \frac{\phi}{2}$

Vertical Force $F_z = rb \int_{\theta_r}^{\theta_f} \{ \tau_x(\theta) \sin\theta + \sigma(\theta) \cos\theta \} d\theta$

Tracked Vehicles



Bekker soil sinkage equation

$$P = k z_0^n = \left(\frac{k_c}{b} + k_\phi \right) z_0^n \Rightarrow z_0 = \left(\frac{P}{k_c/b + k_\phi} \right)^{1/n}$$

Uniform Pressure $\Rightarrow P = \frac{W}{A} = \frac{W}{bl}$ ($A = bl$)

$$\begin{aligned} z_0 &= \left[\frac{W/bl}{k_c/b + k_\phi} \right]^{1/n} \\ \text{Work} &= bl \int_0^{z_0} P dz = bl \int_0^{z_0} \left(\frac{k_c}{b} + k_\phi \right) z^n dz \\ &= bl \left(\frac{k_c}{b} + k_\phi \right) \frac{z_0^{n+1}}{n+1} \end{aligned}$$

Substituting, $\text{Work} = \frac{bl}{(n+1) \left(\frac{k_c}{b} + k_\phi \right)^{1/n}} \left(\frac{W}{bl} \right)^{\frac{n+1}{n}} = R_c l$

$$R_c = \frac{b}{(n+1) \left(\frac{k_c}{b} + k_\phi \right)^{1/n}} \left(\frac{W}{bl} \right)^{\frac{n+1}{n}}$$

$$= \frac{1}{(n+1) b^{1/n} \left(\frac{k_c}{b} + k_\phi \right)^{1/n}} \left(\frac{W}{l} \right)^{\frac{n+1}{n}}$$

$$= \frac{1}{(n+1) (k_c + k_\phi b)^{1/n}} \left(\frac{W}{l} \right)^{\frac{n+1}{n}}$$

motion resistance due to soil compaction by uniformly loaded tread